

Spin of Slater determinant

Spin operators

electron spin

$$\hat{s} = \hat{s}_x \mathbf{x} + \hat{s}_y \mathbf{y} + \hat{s}_z \mathbf{z} \quad \hat{s}^2 = \hat{s}_x^2 + \hat{s}_y^2 + \hat{s}_z^2$$

$$\hat{s}_x |\alpha\rangle = \frac{1}{2} |\beta\rangle \quad \hat{s}_y |\alpha\rangle = \frac{i}{2} |\beta\rangle \quad \hat{s}_z |\alpha\rangle = \frac{1}{2} |\alpha\rangle$$

$$\hat{s}_x |\beta\rangle = \frac{1}{2} |\alpha\rangle \quad \hat{s}_y |\beta\rangle = -\frac{i}{2} |\alpha\rangle \quad \hat{s}_z |\beta\rangle = -\frac{1}{2} |\beta\rangle$$

commutation relations

$$[\hat{S}_x, \hat{S}_y] = \hat{S}_x \hat{S}_y - \hat{S}_y \hat{S}_x = i \hat{S}_z$$

$$[\hat{S}_y, \hat{S}_z] = i \hat{S}_x$$

$$[\hat{S}_z, \hat{S}_x] = i \hat{S}_y$$

Spin of Slater determinant

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$$\hat{s} = \hat{s}_x \mathbf{x} + \hat{s}_y \mathbf{y} + \hat{s}_z \mathbf{z} \quad \hat{s}^2 = \hat{s}_x^2 + \hat{s}_y^2 + \hat{s}_z^2$$

$$\hat{s}_x |\alpha\rangle = \frac{1}{2} |\beta\rangle \quad \hat{s}_y |\alpha\rangle = \frac{i}{2} |\beta\rangle \quad \hat{s}_z |\alpha\rangle = \frac{1}{2} |\alpha\rangle$$

$$\hat{s}_x |\beta\rangle = \frac{1}{2} |\alpha\rangle \quad \hat{s}_y |\beta\rangle = -\frac{i}{2} |\alpha\rangle \quad \hat{s}_z |\beta\rangle = -\frac{1}{2} |\beta\rangle$$

step operators

$$\hat{s}_+ = \hat{s}_x + i\hat{s}_y \quad \hat{s}_+ |\alpha\rangle = 0 \quad \hat{s}_+ |\beta\rangle = |\alpha\rangle$$

$$\hat{s}_- = \hat{s}_x - i\hat{s}_y \quad \hat{s}_- |\alpha\rangle = |\beta\rangle \quad \hat{s}_- |\beta\rangle = 0$$

$$\hat{s}^2 = \hat{s}_+ \hat{s}_- - \hat{s}_z + \hat{s}_z^2$$

$$\hat{s}^2 = \hat{s}_- \hat{s}_+ + \hat{s}_z + \hat{s}_z^2$$

Spin of Slater determinant

many electron spin operators

$$\hat{S}^2 = \hat{S} \cdot \hat{S} = \sum_k^N \sum_l^N \hat{s}(k) \cdot \hat{s}(l) \quad \hat{S}_z = \sum_k^N \hat{s}_z(k)$$

$$\hat{S}_+ = \sum_k^N \hat{s}_+(k) \quad \hat{S}_- = \sum_k^N \hat{s}_-(k)$$

Spin eigenfunctions

$$\hat{S}_z |\phi_1 \phi_2 \dots \phi_3\rangle = \frac{1}{2} (N^\alpha - N^\beta) |\phi_1 \phi_2 \dots \phi_3\rangle$$

closed shell

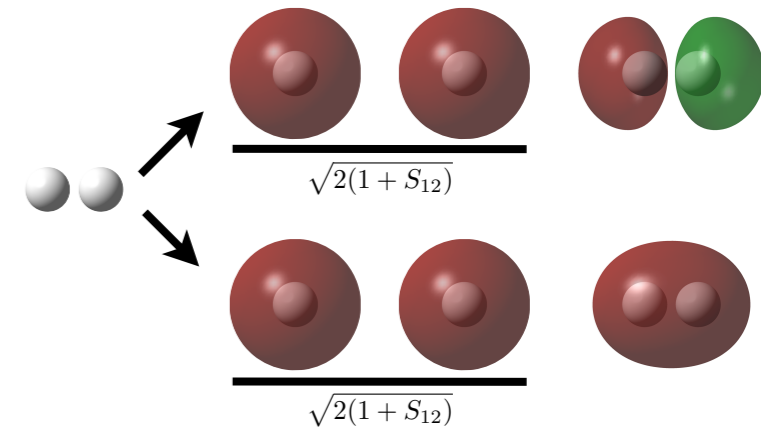
$$\hat{S}^2 |\phi_1 \phi_2 \dots \phi_3\rangle = S(S + 1) |\phi_1 \phi_2 \dots \phi_3\rangle$$

open shell?

configuration state functions

Spin of Slater determinant

minimal basis hydrogen molecule



restricted Hartree-Fock wave function

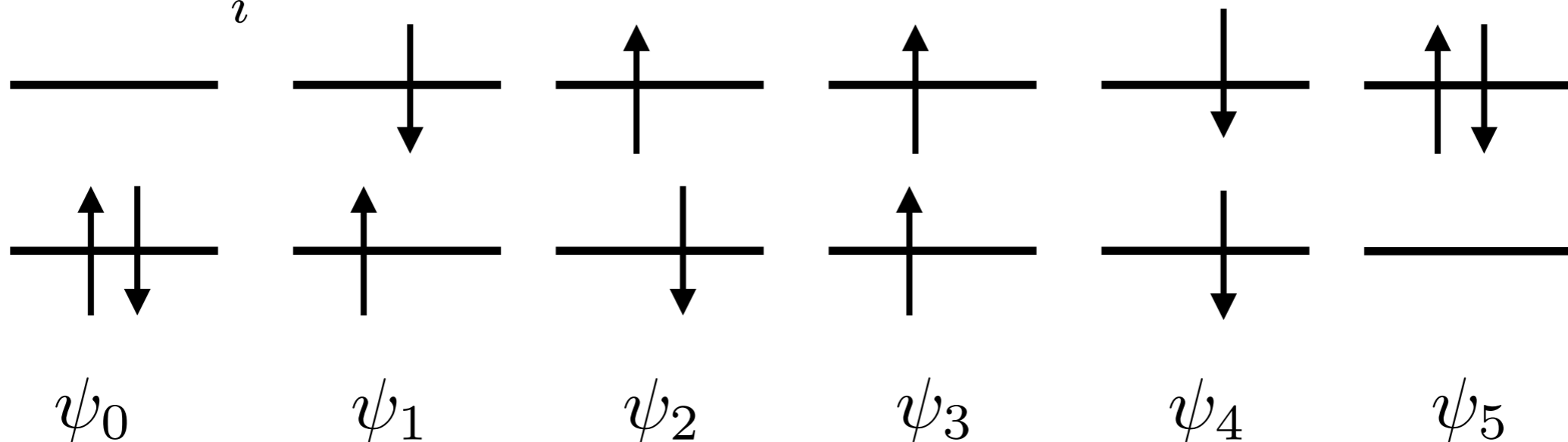
$$\psi_0(\mathbf{x}_1, \mathbf{x}_2) = \frac{1}{\sqrt{2}} [\phi_1(\mathbf{x}_1)\phi_2(\mathbf{x}_2) - \phi_2(\mathbf{x}_1)\phi_1(\mathbf{x}_2)]$$

$$\phi_1(\mathbf{x}) = \varphi_1(\mathbf{r})\alpha(s) \quad \phi_3(\mathbf{x}) = \varphi_2(\mathbf{r})\alpha(s)$$

$$\phi_2(\mathbf{x}) = \varphi_1(\mathbf{r})\beta(s) \quad \phi_4(\mathbf{x}) = \varphi_2(\mathbf{r})\beta(s)$$

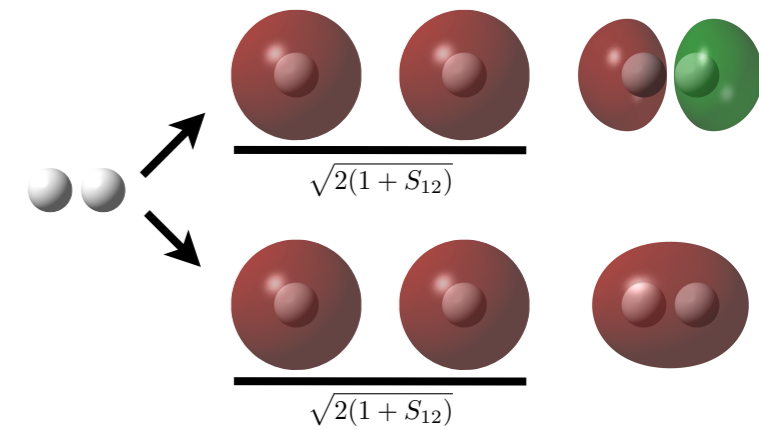
configuration interaction

$$\Psi = \sum_i \psi_i \quad \binom{4}{2} = 6$$

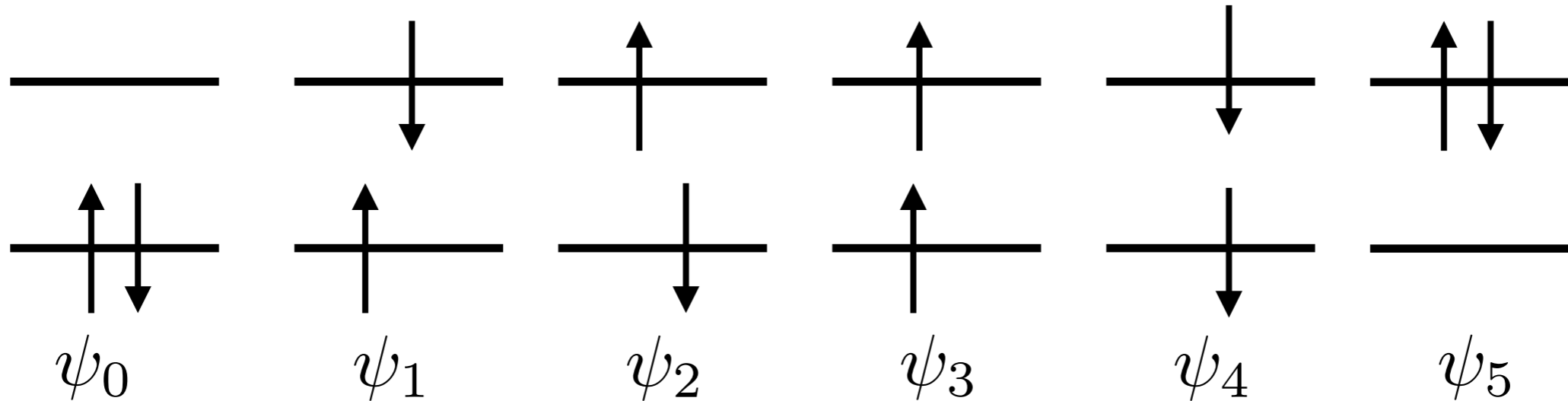


Spin of Slater determinant

minimal basis hydrogen molecule



configurations

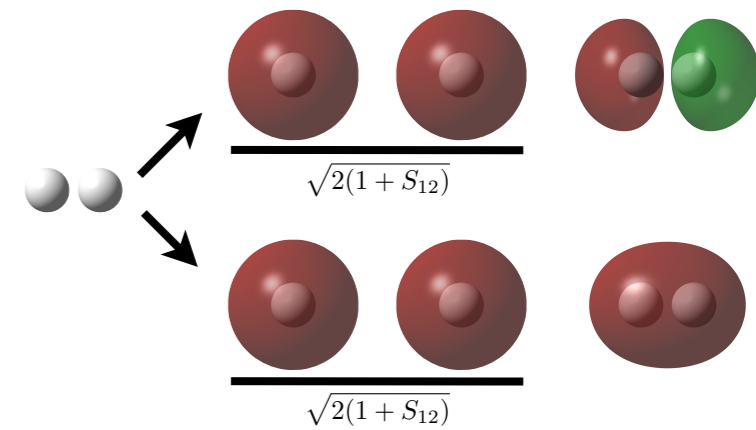


total spin

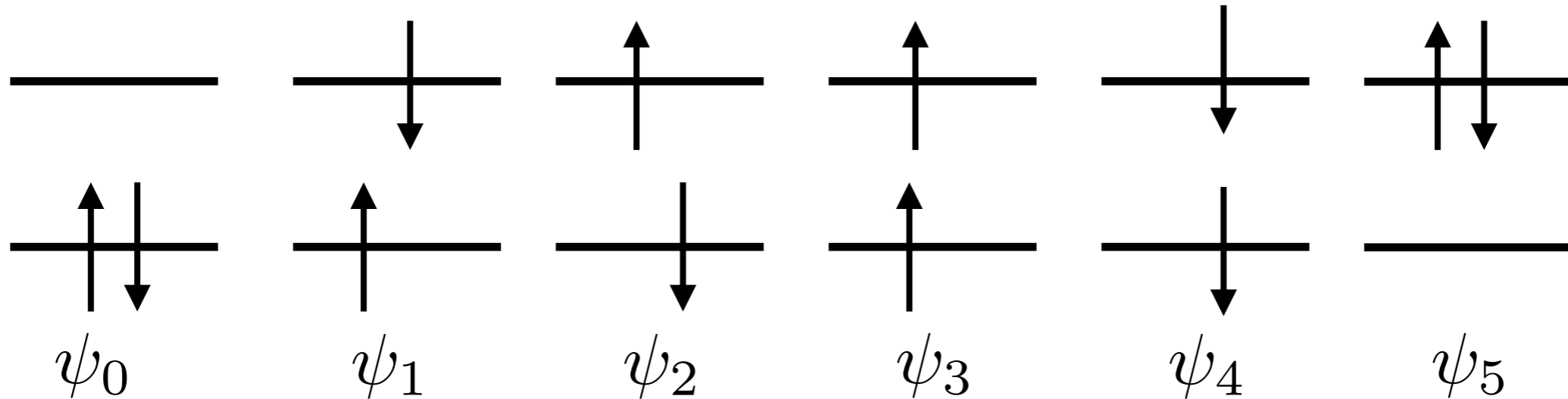
$$\begin{aligned}
 S^2\psi_0 &= (S_-S_+ + S_z + S_z^2)\psi_0 = S(S+1)\psi_0 \\
 &= \frac{1}{\sqrt{2}}[(s_-(1)s_+(1) + s_-(1)s_+(2) + s_-(2)s_+(1) + s_-(2)s_+(2))(\phi_1(1)\phi_2(2) - \phi_2(1)\phi_1(2)) \\
 &\quad + s_z(1)(\phi_1(1)\phi_2(2) - \phi_2(1)\phi_1(2)) + s_z(2)(\phi_1(1)\phi_2(2) - \phi_2(1)\phi_1(2)) \\
 &\quad + (s_z(1)s_z(1) + s_z(1)s_z(2) + s_z(2)s_z(1) + s_z(2)s_z(2))(\phi_1(1)\phi_2(2) - \phi_2(1)\phi_1(2))]
 \end{aligned}$$

Spin of Slater determinant

minimal basis hydrogen molecule



configurations

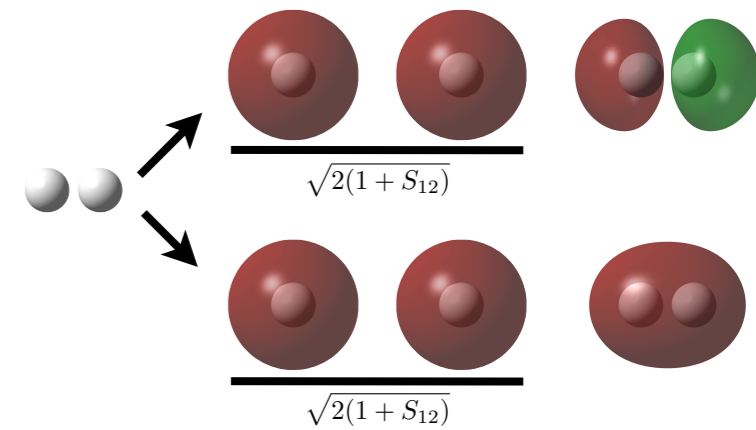


total spin

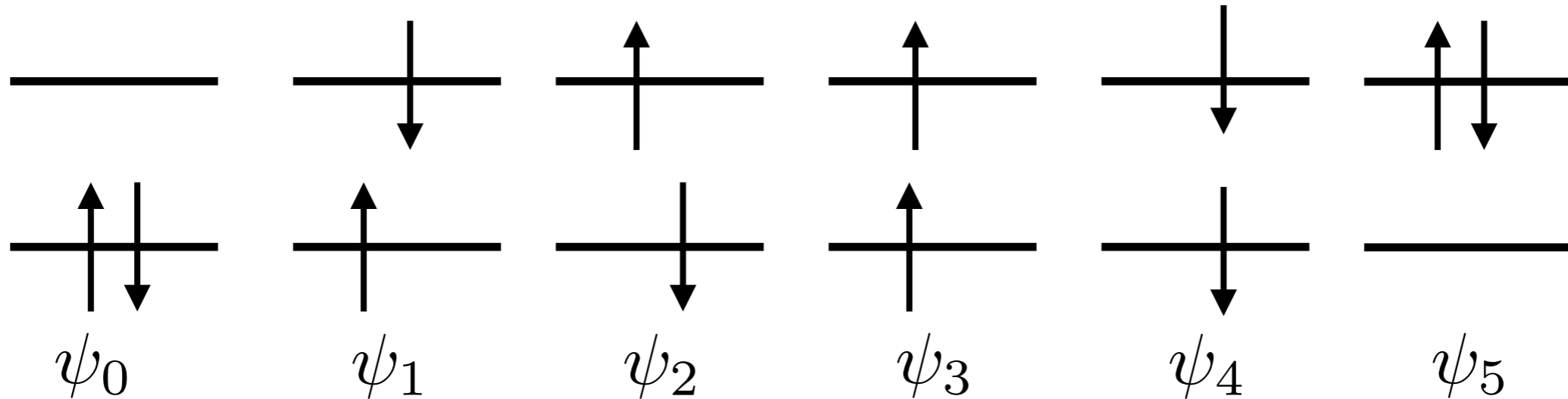
$$\begin{aligned}
 S^2\psi_0 &= \frac{1}{\sqrt{2}}[(s_-(1)s_+(1) + s_-(1)s_+(2) + s_-(2)s_+(1) + s_-(2)s_+(2))\varphi_1(1)\varphi_1(2)(\alpha(1)\beta(2) - \beta(1)\alpha(2)) \\
 &+ s_z(1)\varphi_1(1)\varphi_1(2)(\alpha(1)\beta(2) - \beta(1)\alpha(2)) \\
 &+ s_z(2)\varphi_1(1)\varphi_1(2)(\alpha(1)\beta(2) - \beta(1)\alpha(2)) \\
 &+ (s_z(1)s_z(1) + s_z(1)s_z(2) + s_z(2)s_z(1) + s_z(2)s_z(2))\varphi_1(1)\varphi_1(2)(\alpha(1)\beta(2) - \beta(1)\alpha(2))]
 \end{aligned}$$

Spin of Slater determinant

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configurations

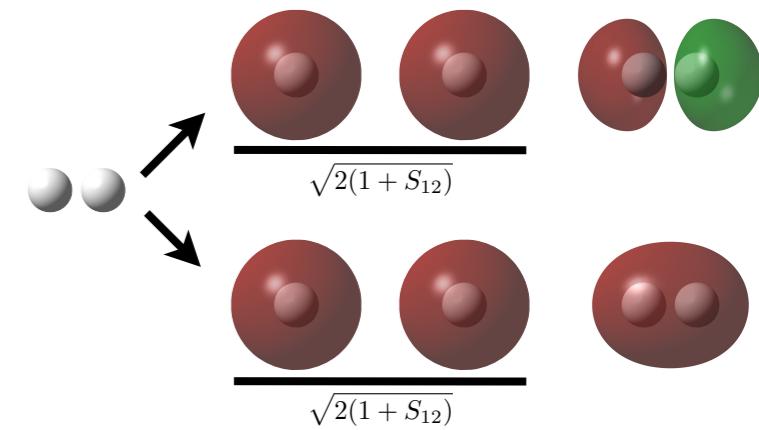


total spin

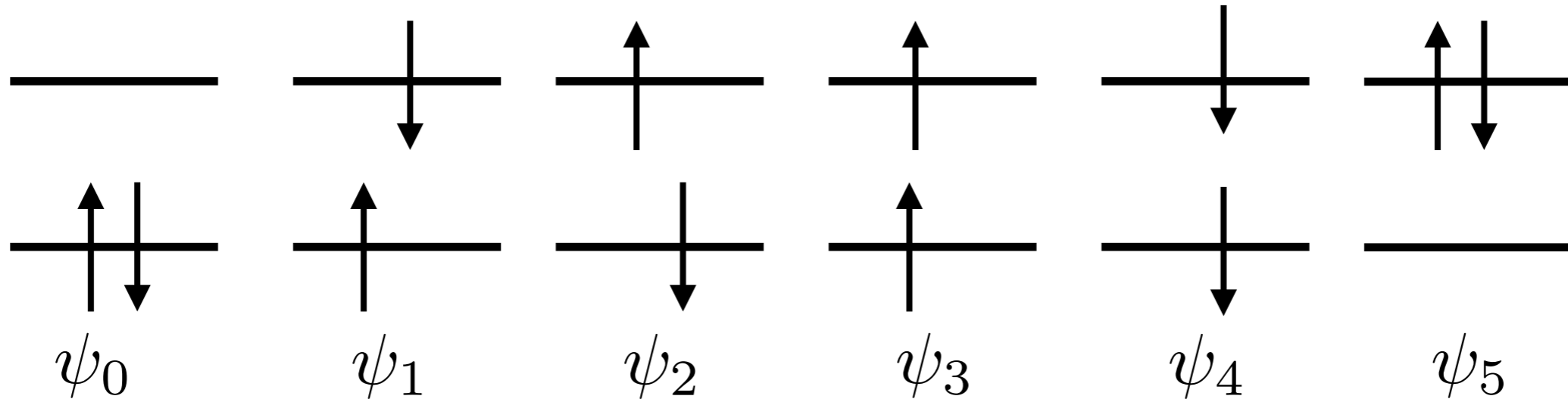
$$\begin{aligned}
 S^2\psi_0 &= \frac{1}{\sqrt{2}}\varphi_1(1)\varphi_1(2)[(s_-(1)s_+(1) + s_-(1)s_+(2) + s_-(2)s_+(1) + s_-(2)s_+(2))(\alpha(1)\beta(2) - \beta(1)\alpha(2)) \\
 &\quad + s_z(1)(\alpha(1)\beta(2) - \beta(1)\alpha(2)) \\
 &\quad + s_z(2)(\alpha(1)\beta(2) - \beta(1)\alpha(2)) \\
 &\quad + (s_z(1)s_z(1) + s_z(1)s_z(2) + s_z(2)s_z(1) + s_z(2)s_z(2))(\alpha(1)\beta(2) - \beta(1)\alpha(2))]
 \end{aligned}$$

Spin of Slater determinant

minimal basis hydrogen molecule



configurations



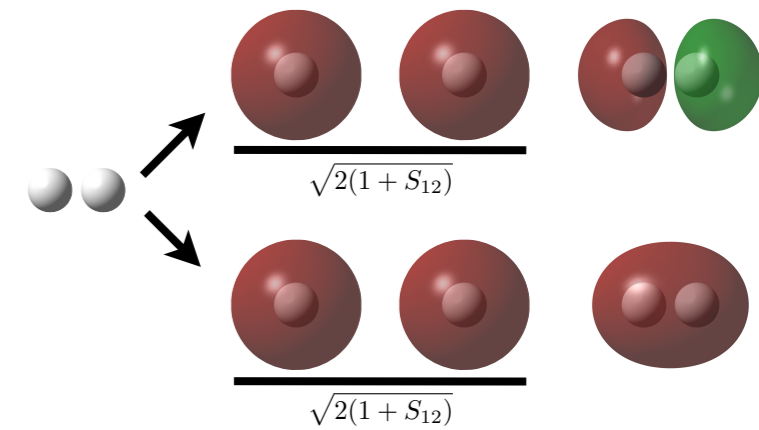
total spin

$$\begin{aligned}
 S^2 \psi_0 &= \frac{1}{\sqrt{2}} \varphi_1(1) \varphi_1(2) [-\beta(1)\alpha(2) + \beta(1)\alpha(2) - \alpha(1)\beta(2) + \alpha(1)\beta(2) \\
 &\quad + \frac{1}{2}\alpha(1)\beta(2) + \frac{1}{2}\beta(1)\alpha(2) \\
 &\quad - \frac{1}{2}\alpha(1)\beta(2) - \frac{1}{2}\beta(1)\alpha(2) \\
 &\quad + (\frac{1}{4} - \frac{1}{4} - \frac{1}{4} + \frac{1}{4} - \frac{1}{4} + \frac{1}{4} + \frac{1}{4} - \frac{1}{4})(\alpha(1)\beta(2) - \beta(1)\alpha(2))] \\
 &= 0
 \end{aligned}$$

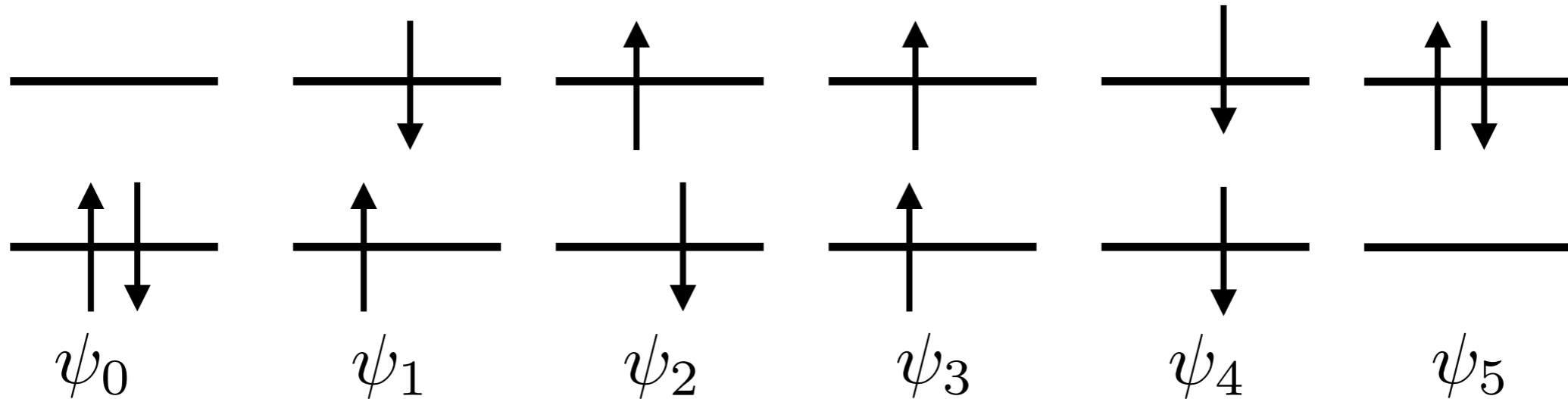
singlet

Spin of Slater determinant

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configurations



total spin

$$\psi_3 = (\varphi_1(1)\varphi_2(2) - \varphi_2(1)\varphi_1(2))\alpha(1)\alpha(2)$$

$$S^2\psi_3 = (S_-S_+ + S_z + S_z^2)\psi_3 = S(S+1)\psi_3$$

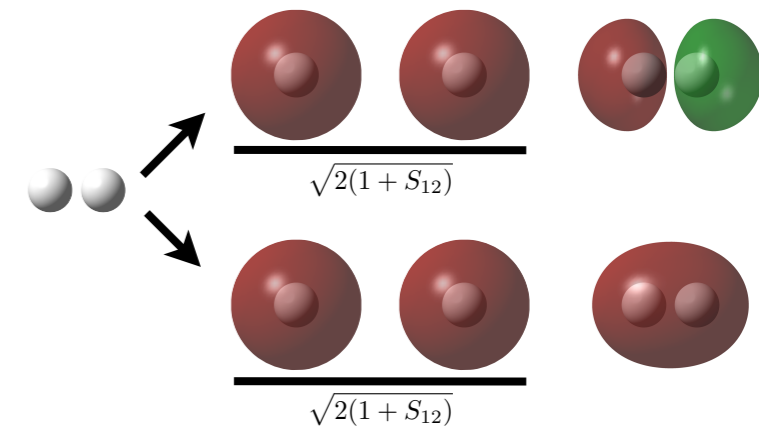
$$= \frac{1}{\sqrt{2}}(\varphi_1(1)\varphi_2(2) - \varphi_2(1)\varphi_1(2))[(s_-(1)s_+(1) + s_-(1)s_+(2) + s_-(2)s_+(1) + s_-(2)s_+(2))\alpha(1)\alpha(2)$$

$$+s_z(1)\alpha(1)\alpha(2) + s_z(2)\alpha(1)\alpha(2)$$

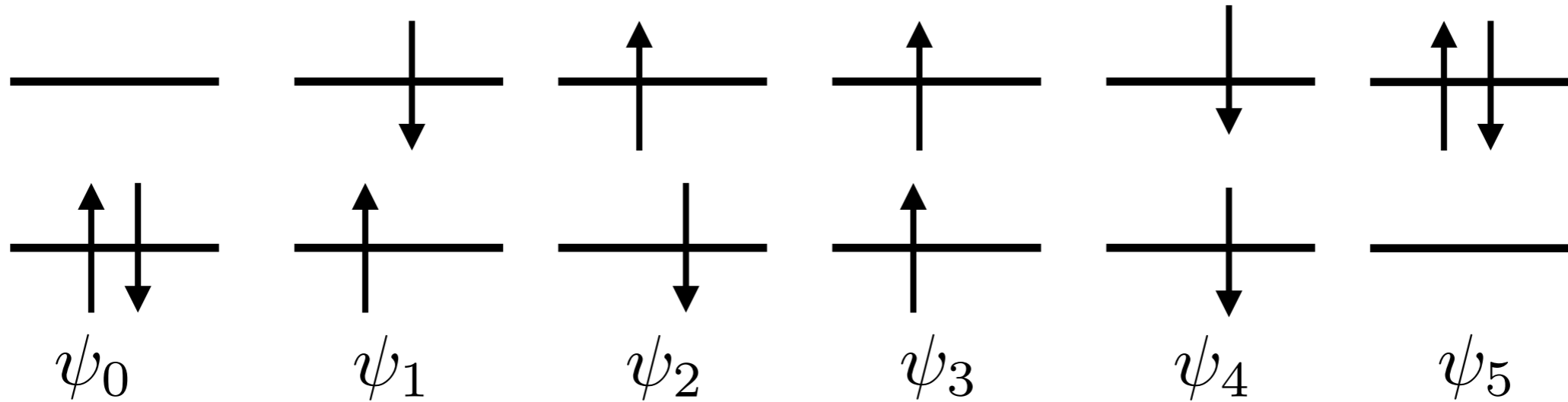
$$+(s_z(1)s_z(1) + s_z(1)s_z(2) + s_z(2)s_z(1) + s_z(2)s_z(2))\alpha(1)\alpha(2)]$$

Spin of Slater determinant

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configurations



total spin

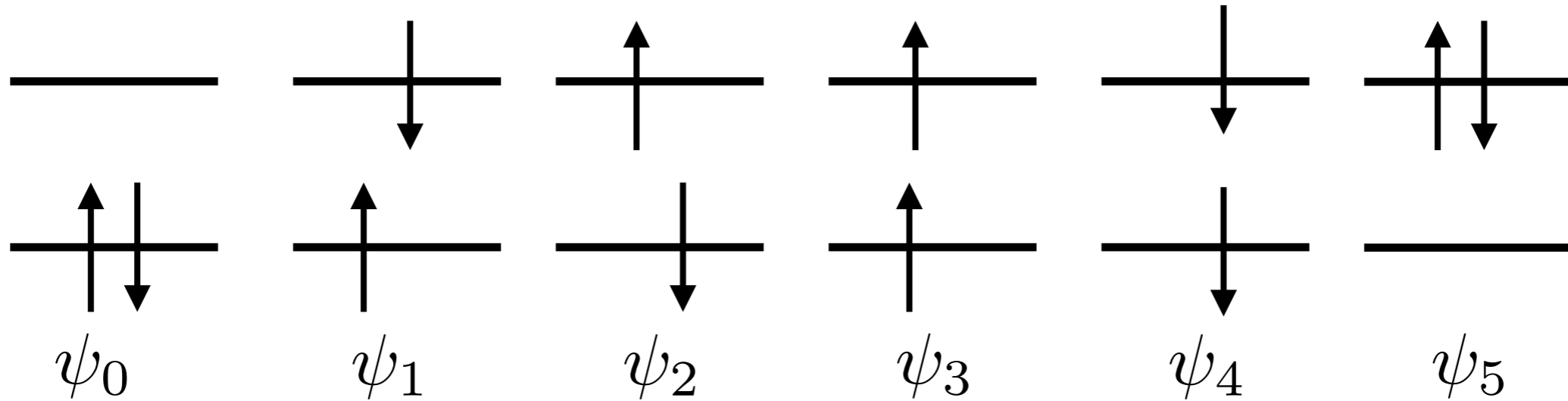
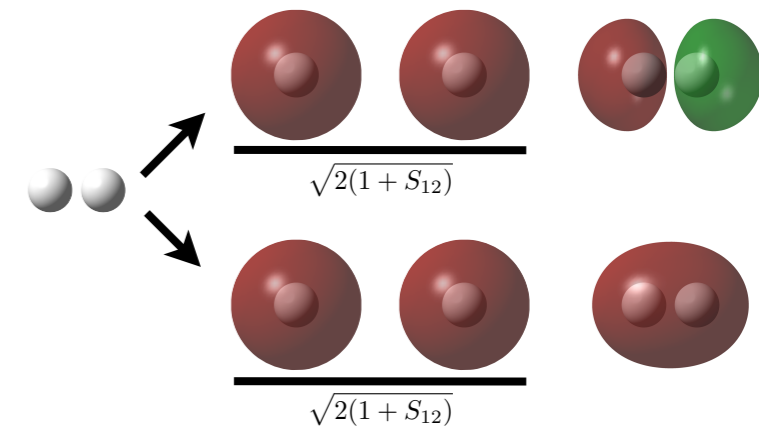
$$\begin{aligned}
 S^2\psi_3 &= \frac{1}{\sqrt{2}}(\varphi_1(1)\varphi_2(2) - \varphi_2(1)\varphi_1(2)) [\\
 &\quad (0 + 0 + 0 + 0)\alpha(1)\alpha(2) \\
 &\quad + (\frac{1}{2} + \frac{1}{2})\alpha(1)\alpha(2) \\
 &\quad + (\frac{1}{4} + \frac{1}{4} + \frac{1}{4} + \frac{1}{4})\alpha(1)\alpha(2)] \\
 &= 2\frac{1}{\sqrt{2}}(\varphi_1(1)\varphi_2(2) - \varphi_2(1)\varphi_1(2))\alpha(1)\alpha(2)
 \end{aligned}$$

triplet

Spin of Slater determinant

minimal basis hydrogen molecule

configurations



total spin

$$\psi_4 = (\varphi_1(1)\varphi_2(2) - \varphi_2(1)\varphi_1(2))\beta(1)\beta(2)$$

$$S^2\psi_4 = (S_-S_+ + S_z + S_z^2)\psi_4 = S(S+1)\psi_4$$

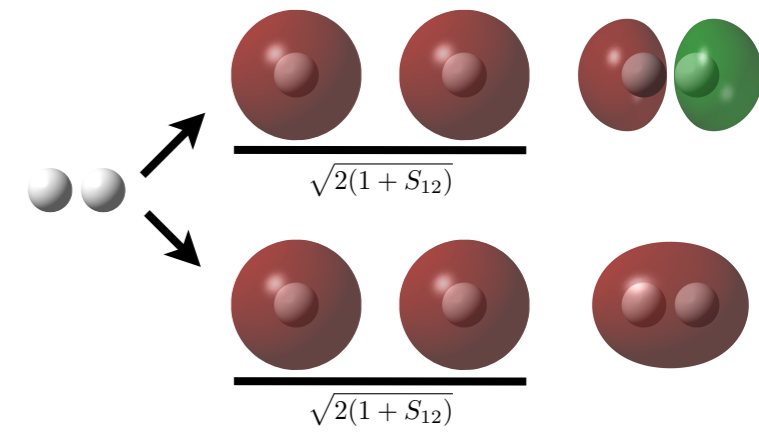
$$= \frac{1}{\sqrt{2}}(\varphi_1(1)\varphi_2(2) - \varphi_2(1)\varphi_1(2))[(s_-(1)s_+(1) + s_-(1)s_+(2) + s_-(2)s_+(1) + s_-(2)s_+(2))\beta(1)\beta(2)$$

$$+s_z(1)\beta(1)\beta(2) + s_z(2)\beta(1)\beta(2)$$

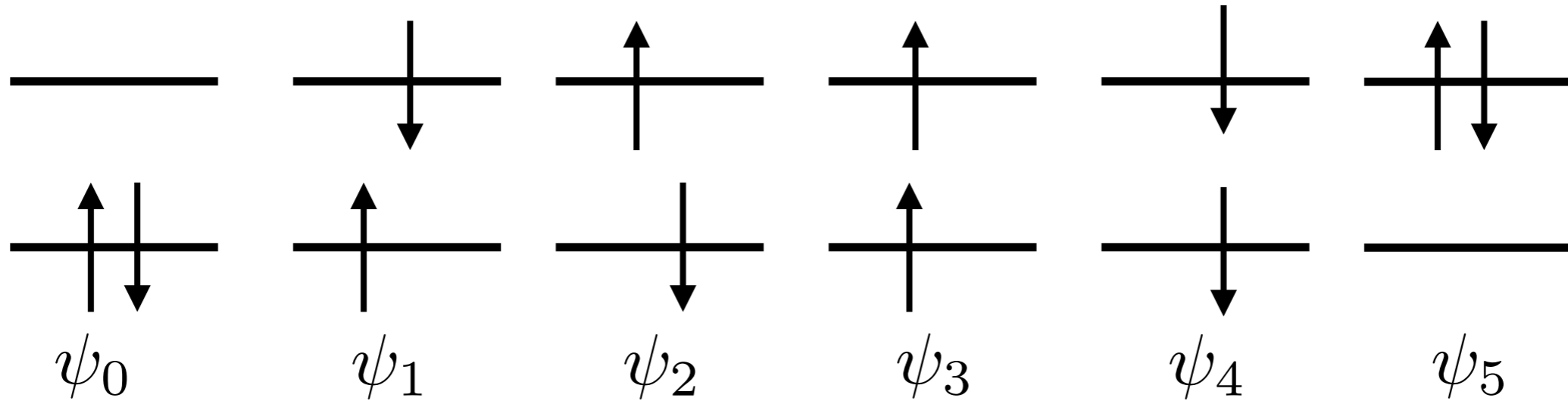
$$+(s_z(1)s_z(1) + s_z(1)s_z(2) + s_z(2)s_z(1) + s_z(2)s_z(2))\beta(1)\beta(2)]$$

Spin of Slater determinant

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configurations



total spin

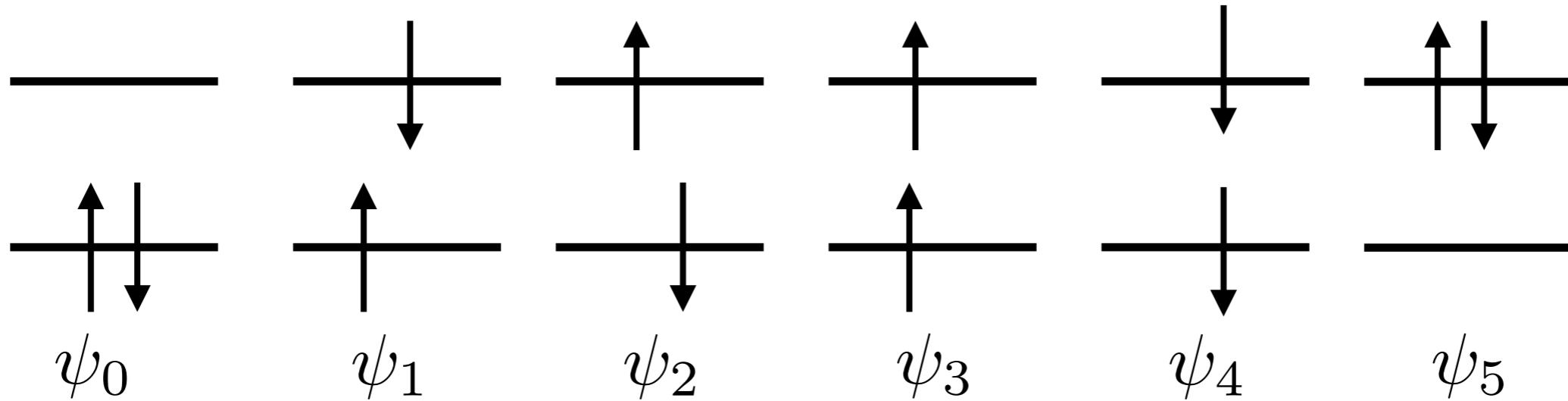
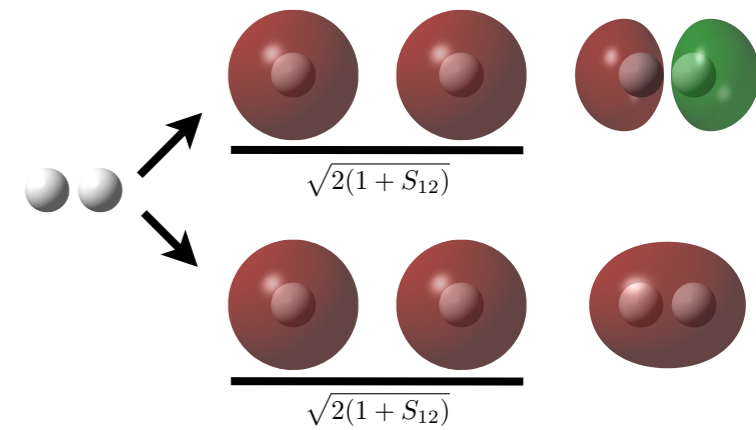
$$\begin{aligned}
 S^2\psi_4 &= \frac{1}{\sqrt{2}}(\varphi_1(1)\varphi_2(2) - \varphi_2(1)\varphi_1(2)) [\\
 &\quad (1 + 0 + 0 + 1)\beta(1)\beta(2) \\
 &\quad + (-\frac{1}{2} - \frac{1}{2})\beta(1)\beta(2) \\
 &\quad + (\frac{1}{4} + \frac{1}{4} + \frac{1}{4} + \frac{1}{4})\beta(1)\beta(2)] \\
 &= 2\frac{1}{\sqrt{2}}(\varphi_1(1)\varphi_2(2) - \varphi_2(1)\varphi_1(2))\beta(1)\beta(2)
 \end{aligned}$$

triplet

Spin of Slater determinant

minimal basis hydrogen molecule

configurations

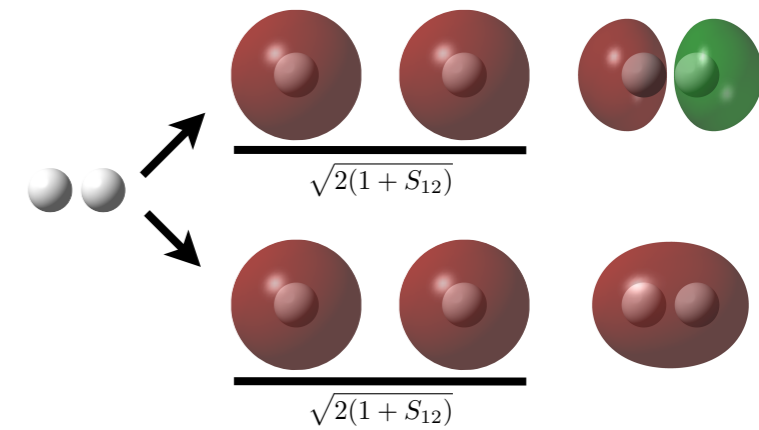


total spin

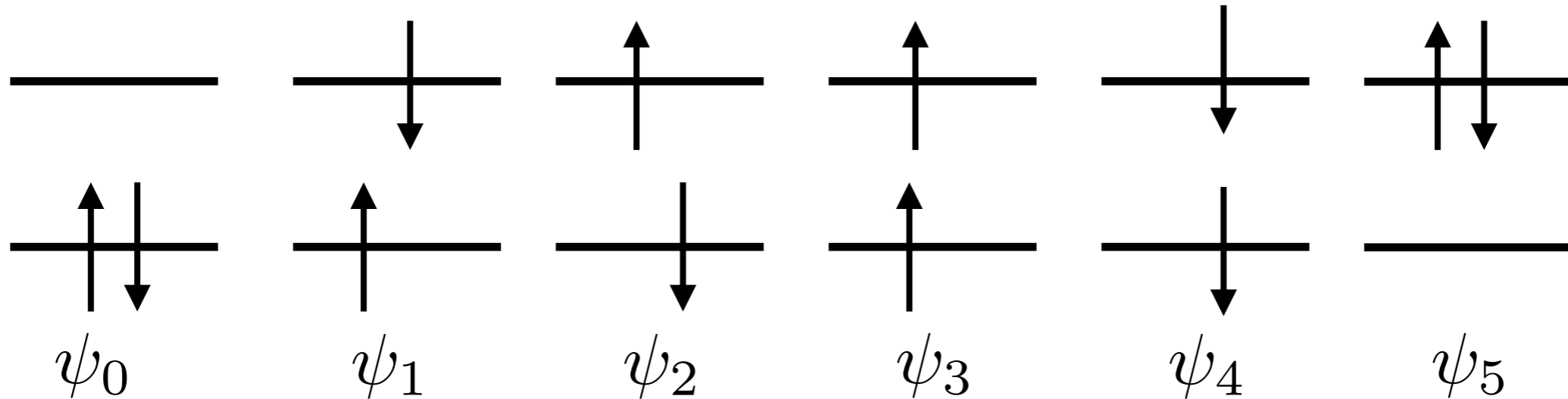
$$\psi_1 = \frac{1}{\sqrt{2}} [\varphi_1(1)\alpha(1)\varphi_2\beta(2) - \varphi_2(1)\beta(2)\varphi_1(2)\alpha(2)]$$

Spin of Slater determinant

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configurations

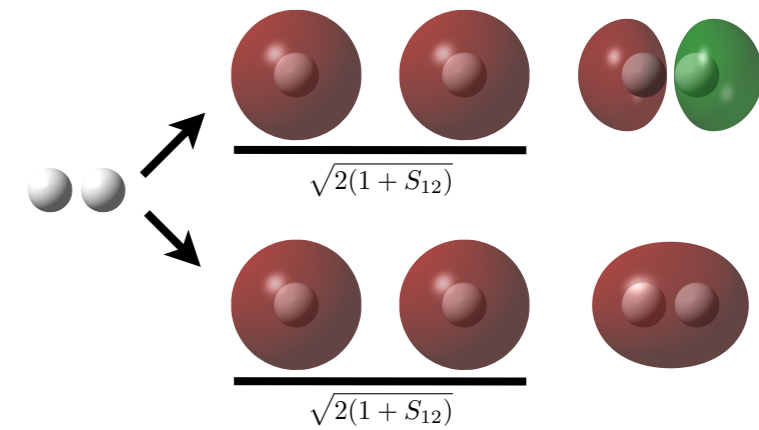


total spin

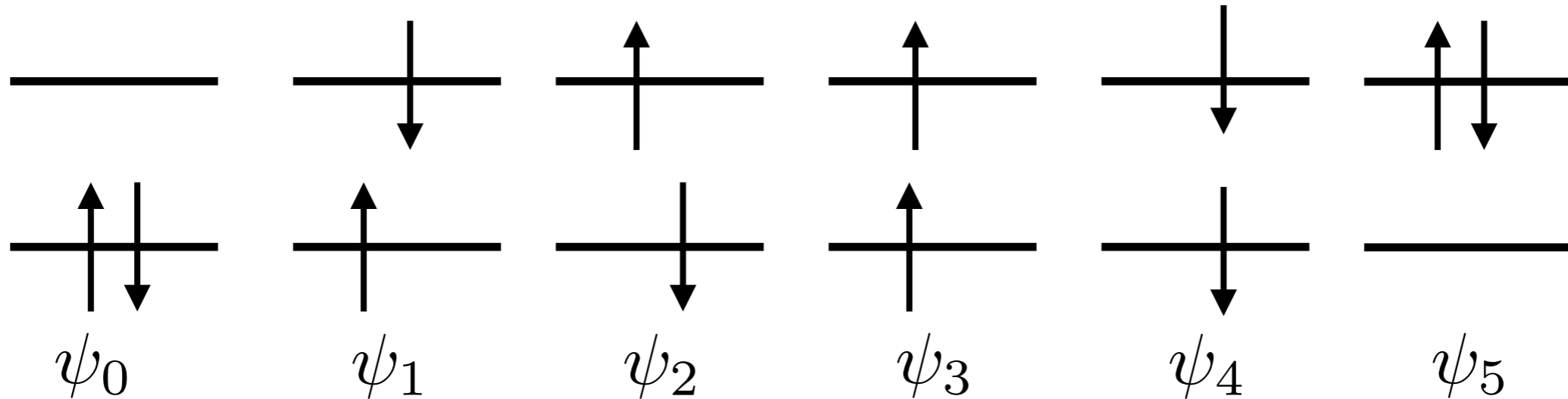
$$\begin{aligned}
 S^2\psi_1 = & \frac{1}{\sqrt{2}}[(s_-(1)s_+(1) + s_-(1)s_+(2) + s_-(2)s_+(1) + s_-(2)s_+(2))\varphi_1(1)\varphi_2(2)\alpha(1)\beta(2) \\
 & - (s_-(1)s_+(1) + s_-(1)s_+(2) + s_-(2)s_+(1) + s_-(2)s_+(2))\varphi_2(1)\varphi_1(2)\beta(1)\alpha(2) \\
 & + s_z(1)\varphi_1(1)\varphi_2(2)\alpha(1)\beta(2) - s_z(1)\varphi_2(1)\varphi_1(2)\beta(1)\alpha(2) \\
 & + s_z(2)\varphi_1(1)\varphi_2(2)\alpha(1)\beta(2) - s_z(2)\varphi_2(1)\varphi_1(2)\beta(1)\alpha(2) \\
 & + (s_z(1)s_z(1) + s_z(1)s_z(2) + s_z(2)s_z(1) + s_z(2)s_z(2))\varphi_1(1)\varphi_2(2)\alpha(1)\beta(2) \\
 & - (s_z(1)s_z(1) + s_z(1)s_z(2) + s_z(2)s_z(1) + s_z(2)s_z(2))\varphi_2(1)\varphi_1(2)\beta(1)\alpha(2)
 \end{aligned}$$

Spin of Slater determinant

minimal basis hydrogen molecule



configurations



total spin

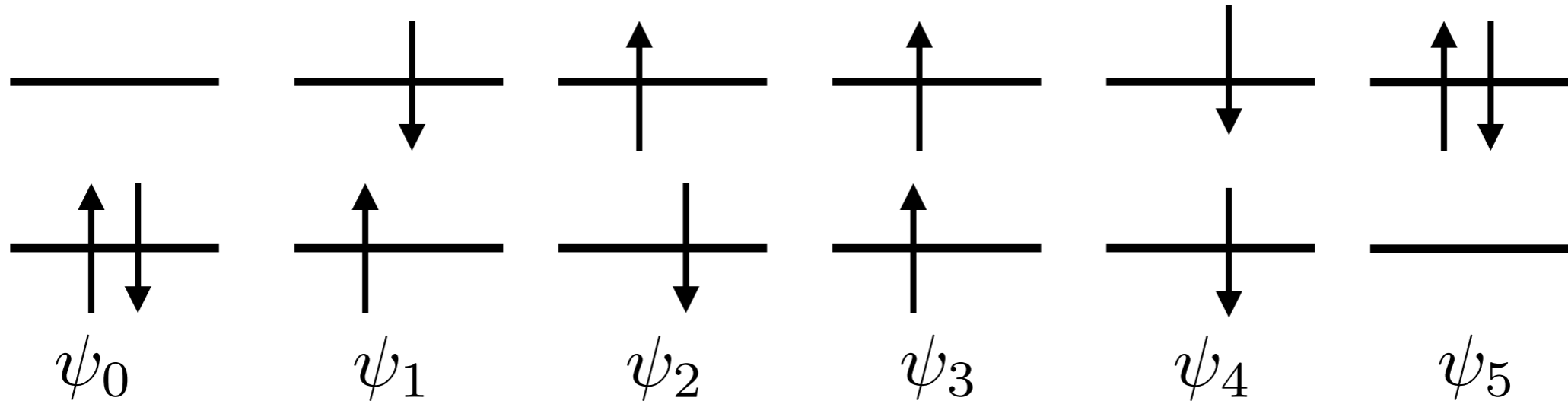
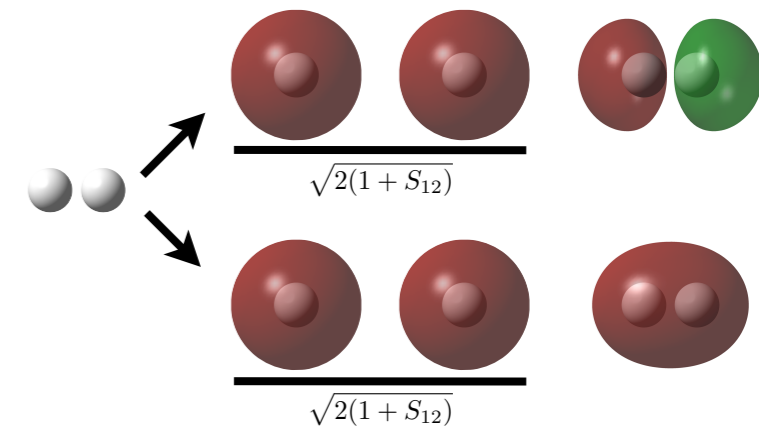
$$\begin{aligned} S^2\psi_1 &= \frac{1}{\sqrt{2}}[\varphi_1(1)\alpha(1)\varphi_2(2)\beta(2) - \varphi_2(1)\beta(1)\varphi_1(2)\alpha(2) \\ &\quad + \varphi_1(1)\beta(1)\varphi_2(2)\alpha(2) - \varphi_2(1)\alpha(1)\varphi_1(2)\beta(2)] \\ &= \psi_1 + \psi_2 \\ &\neq S(S+1)\psi_1 \end{aligned}$$

no eigenfunction of total spin operator!

Spin of Slater determinant

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configurations



spin-adapted linear combinations:

$$\psi_A = \psi_1 + \psi_2$$

$$= (\varphi_1(1)\varphi_2(1) - \varphi_2(1)\varphi_1(2))(\alpha(1)\beta(2) + \beta(1)\alpha(2))$$

$$\psi_B = \psi_1 - \psi_2$$

$$= (\varphi_1(1)\varphi_2(1) + \varphi_2(1)\varphi_1(2))(\alpha(1)\beta(2) - \beta(1)\alpha(2))$$

homework: show that both are eigenfunction of total spin operator.

Hartree-Fock Theory for n electrons

Expectations values for one and two electron operators

Hartree product (no spin)

$$\langle \hat{O}_1 \rangle = \sum_a \int \phi_a^*(\mathbf{x}_1) \hat{o}(\mathbf{r}_1) \phi_a(\mathbf{r}_1) d\mathbf{r}_1$$

$$\langle \hat{O}_2 \rangle = \frac{1}{2} \sum_a \sum_b \int \int \phi_a^*(\mathbf{r}_1) \phi_b^*(\mathbf{r}_2) \hat{o}(\mathbf{r}_1, \mathbf{r}_2) \phi_a(\mathbf{r}_1) \phi_b(\mathbf{r}_2) d\mathbf{r}_1 d\mathbf{r}_2$$

Slater determinant (spin, Pauli principle)

$$\langle \hat{O}_1 \rangle = \sum_a \int \phi_a^*(\mathbf{x}_1) \hat{o}(\mathbf{r}_1) \phi_a(\mathbf{x}_1) d\mathbf{x}_1$$

$$\begin{aligned} \langle \hat{O}_2 \rangle &= \frac{1}{2} \sum_a \sum_b \int \int \phi_a^*(\mathbf{x}_1) \phi_b^*(\mathbf{x}_2) \hat{o}(\mathbf{r}_1, \mathbf{r}_2) \phi_a(\mathbf{x}_1) \phi_b(\mathbf{x}_2) d\mathbf{x}_1 d\mathbf{x}_2 \\ &\quad - \frac{1}{2} \sum_a \sum_b \int \int \phi_a^*(\mathbf{x}_1) \phi_b^*(\mathbf{x}_2) \hat{o}(\mathbf{r}_1, \mathbf{r}_2) \phi_b(\mathbf{x}_1) \phi_a(\mathbf{x}_2) d\mathbf{x}_1 d\mathbf{x}_2 \end{aligned}$$

$$\langle \hat{O}_2 \rangle = \frac{1}{2} \sum_a \sum_b \int \int \phi_a^*(\mathbf{x}_1) \phi_b^*(\mathbf{x}_2) \hat{o}(\mathbf{r}_1, \mathbf{r}_2) (1 - \hat{p}_{12}) \phi_a(\mathbf{x}_1) \phi_b(\mathbf{x}_2) d\mathbf{x}_1 d\mathbf{x}_2$$

Hartree-Fock Theory for n electrons

Expectations values for one and two electron operators

Hartree product (no spin)

$$\langle \hat{O}_1 \rangle = \sum_a \int \phi_a^*(\mathbf{x}_1) \hat{o}(\mathbf{r}_1) \phi_a(\mathbf{r}_1) d\mathbf{r}_1$$

$$\langle \hat{O}_2 \rangle = \frac{1}{2} \sum_a \sum_b \int \int \phi_a^*(\mathbf{r}_1) \phi_b^*(\mathbf{r}_2) \hat{o}(\mathbf{r}_1, \mathbf{r}_2) \phi_a(\mathbf{r}_1) \phi_b(\mathbf{r}_2) d\mathbf{r}_1 d\mathbf{r}_2$$

Slater determinant (spin, Pauli principle)

$$\langle \hat{O}_1 \rangle = \sum_a \int \phi_a^*(\mathbf{x}_1) \hat{o}(\mathbf{r}_1) \phi_a(\mathbf{x}_1) d\mathbf{x}_1$$

$$\langle \hat{O}_2 \rangle = \frac{1}{2} \sum_a \sum_b \int \int \phi_a^*(\mathbf{x}_1) \phi_b^*(\mathbf{x}_2) \hat{o}(\mathbf{r}_1, \mathbf{r}_2) \phi_a(\mathbf{x}_1) \phi_b(\mathbf{x}_2) d\mathbf{x}_1 d\mathbf{x}_2$$

$$-\frac{1}{2} \sum_a \sum_b \int \int \phi_a^*(\mathbf{x}_1) \phi_b^*(\mathbf{x}_2) \hat{o}(\mathbf{r}_1, \mathbf{r}_2) \phi_b(\mathbf{x}_1) \phi_a(\mathbf{x}_2) d\mathbf{x}_1 d\mathbf{x}_2$$

$$\langle \hat{O}_2 \rangle = \frac{1}{2} \sum_a \sum_b \int \int \phi_a^*(\mathbf{x}_1) \phi_b^*(\mathbf{x}_2) \hat{o}(\mathbf{r}_1, \mathbf{r}_2) (1 - \hat{p}_{12}) \phi_a(\mathbf{x}_1) \phi_b(\mathbf{x}_2) d\mathbf{x}_1 d\mathbf{x}_2$$

Hartree-Fock Theory for n electrons

task: finding the best orbitals that minimize total energy

Variational principle

$$E \geq E^0$$

minimum

$$\partial E = 0$$

when for all orbitals

$$\varphi_i(\mathbf{x}) = \varphi_i(\mathbf{x}) + \delta\varphi_i(\mathbf{x})$$

constraint

$$\langle \varphi_i(\mathbf{x}) | \varphi_k(\mathbf{x}) \rangle = \delta_{ik}$$

Lagrange multipliers

Hartree-Fock Theory for n electrons

task: finding the best orbitals that minimize total energy

Lagrangian (adding a smart zero, or actually n^2 smart zeros)

$$\mathcal{L} = E - \sum_i \sum_k \epsilon_{ik} [\langle \varphi_i | \varphi_k \rangle - \delta_{ik}]$$

minimum

$$\partial \mathcal{L} = \partial E - \sum_i \sum_k \epsilon_{ik} [\langle \partial \varphi_i | \varphi_k \rangle + \langle \varphi_i | \partial \varphi_k \rangle] = 0$$

rearranging

$$\begin{aligned} \partial E &= \sum_i \langle \partial \varphi_i | \hat{h} | \varphi_i \rangle + \frac{1}{2} \sum_i \sum_k [\langle \partial \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_i \varphi_k \rangle + \langle \varphi_i \partial \varphi_k | \frac{1}{r_{12}} | \varphi_i \varphi_k \rangle \\ &\quad - \langle \partial \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_k \varphi_i \rangle - \langle \varphi_i \partial \varphi_k | \frac{1}{r_{12}} | \varphi_k \varphi_i \rangle] \\ &+ \sum_i \langle \varphi_i | \hat{h} | \partial \varphi_i \rangle + \frac{1}{2} \sum_i \sum_k [\langle \varphi_i \varphi_k | \frac{1}{r_{12}} | \partial \varphi_i \varphi_k \rangle + \langle \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_i \partial \varphi_k \rangle \\ &\quad - \langle \varphi_i \varphi_k | \frac{1}{r_{12}} | \partial \varphi_k \varphi_i \rangle - \langle \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_k \partial \varphi_i \rangle] \end{aligned}$$

Hartree-Fock Theory for n electrons

task: finding the best orbitals that minimize total energy

Lagrangian (adding a smart zero!)

$$\mathcal{L} = E - \sum_i \sum_k \epsilon_{ik} [\langle \varphi_i | \varphi_k \rangle - \delta_{ik}]$$

minimum

$$\partial \mathcal{L} = \partial E - \sum_i \sum_k \epsilon_{ik} [\langle \partial \varphi_i | \varphi_k \rangle + \langle \varphi_i | \partial \varphi_k \rangle] = 0$$

rearranging

$$\begin{aligned} \partial E = \sum_i \langle \partial \varphi_i | \hat{h} | \varphi_i \rangle + \frac{1}{2} \sum_i \sum_k & [\langle \partial \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_i \varphi_k \rangle + \langle \varphi_i \partial \varphi_k | \frac{1}{r_{12}} | \varphi_i \varphi_k \rangle \\ & - \langle \partial \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_k \varphi_i \rangle - \langle \varphi_i \partial \varphi_k | \frac{1}{r_{12}} | \varphi_k \varphi_i \rangle] \end{aligned}$$

$$\begin{aligned} + \sum_i \langle \varphi_i | \hat{h} | \partial \varphi_i \rangle + \frac{1}{2} \sum_i \sum_k & [\langle \varphi_i \varphi_k | \frac{1}{r_{12}} | \partial \varphi_i \varphi_k \rangle + \langle \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_i \partial \varphi_k \rangle \\ & - \langle \varphi_i \varphi_k | \frac{1}{r_{12}} | \partial \varphi_k \varphi_i \rangle - \langle \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_k \partial \varphi_i \rangle] \end{aligned}$$

complex conjugate

Hartree-Fock Theory for n electrons

task: finding the best orbitals that minimize total energy

Lagrangian (adding a smart zero!)

$$\mathcal{L} = E - \sum_i \sum_k \epsilon_{ik} [\langle \varphi_i | \varphi_k \rangle - \delta_{ik}]$$

minimum

$$\partial \mathcal{L} = \partial E - \sum_i \sum_k \epsilon_{ik} [\langle \partial \varphi_i | \varphi_k \rangle + \langle \varphi_i | \partial \varphi_k \rangle] = 0$$

rearranging

$$\partial E = \sum_i \langle \partial \varphi_i | \hat{h} | \varphi_i \rangle + \frac{1}{2} \sum_i \sum_k$$

changing variables

$$[\langle \partial \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_i \varphi_k \rangle + \langle \varphi_i \partial \varphi_k | \frac{1}{r_{12}} | \varphi_i \varphi_k \rangle - \langle \partial \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_k \varphi_i \rangle - \langle \varphi_i \partial \varphi_k | \frac{1}{r_{12}} | \varphi_k \varphi_i \rangle]$$

$$+ \sum_i \langle \varphi_i | \hat{h} | \partial \varphi_i \rangle + \frac{1}{2} \sum_i \sum_k [\langle \varphi_i \varphi_k | \frac{1}{r_{12}} | \partial \varphi_i \varphi_k \rangle + \langle \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_i \partial \varphi_k \rangle - \langle \varphi_i \varphi_k | \frac{1}{r_{12}} | \partial \varphi_k \varphi_i \rangle - \langle \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_k \partial \varphi_i \rangle]$$

complex conjugate

Hartree-Fock Theory for n electrons

task: finding the best orbitals that minimize total energy

Lagrangian (adding a smart zero!)

$$\mathcal{L} = E - \sum_i \sum_k \epsilon_{ik} [\langle \varphi_i | \varphi_k \rangle - \delta_{ik}]$$

minimum

$$\partial \mathcal{L} = \partial E - \sum_i \sum_k \epsilon_{ik} [\langle \partial \varphi_i | \varphi_k \rangle + \langle \varphi_i | \partial \varphi_k \rangle] = 0$$

rearranging

$$\begin{aligned} \partial E &= \sum_i \langle \partial \varphi_i | \hat{h} | \varphi_i \rangle + \sum_i \sum_k [\langle \partial \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_i \varphi_k \rangle - \langle \partial \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_k \varphi_i \rangle] \\ &+ \text{c.c} \end{aligned}$$

Hartree-Fock Theory for n electrons

task: finding the best orbitals that minimize total energy

rearranging

$$\begin{aligned} \partial E = & \sum_i \langle \partial \varphi_i | \hat{h} | \varphi_i \rangle + \sum_i \sum_k \langle \partial \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_i \varphi_k \rangle && \text{coulomb} \\ & - \sum_i \sum_k \langle \partial \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_k \varphi_i \rangle \\ & + \text{c.c.} && \text{exchange} \end{aligned}$$

for each orbital:

$$\langle \partial \varphi_i | \hat{h} | \varphi_i \rangle + \langle \partial \varphi_i | \sum_k \frac{|\varphi_k|^2}{r_{12}} | \varphi_i \rangle - \langle \partial \varphi_i | \sum_k \langle \varphi_k | \frac{1}{r_{12}} | \varphi_i \rangle \varphi_k \rangle$$

Hartree-Fock Theory for n electrons

task: finding the best orbitals that minimize total energy

rearranging a bit more

$$\langle \partial\varphi_i | \hat{h} | \varphi_i \rangle + \langle \partial\varphi_i | \sum_k \frac{|\varphi_k|^2}{r_{12}} | \varphi_i \rangle - \langle \partial\varphi_i | \sum_k \langle \varphi_k | \frac{1}{r_{12}} | \varphi_i \rangle \varphi_k \rangle$$

local coulomb operator

$$\hat{J}_k | \varphi_i(1) \rangle = \frac{|\varphi_k|^2}{r_{12}} | \varphi_i(1) \rangle$$

non-local exchange operator

$$\hat{K}_k | \varphi_i(1) \rangle = \langle \varphi_k(2) | \frac{1}{r_{12}} | \varphi_i(2) \rangle | \varphi_k(1) \rangle$$

Fock operator

$$\hat{f} = \hat{h} + \sum_k \left[\hat{J}_k - \hat{K}_k \right]$$

Hartree-Fock Theory for n electrons

task: finding the best orbitals that minimize total energy

going back to the Lagrangian

$$\partial \mathcal{L} = \partial E - \sum_i \sum_k \epsilon_{ik} [\langle \partial \varphi_i | \varphi_k \rangle + \langle \varphi_i | \partial \varphi_k \rangle] = 0$$

thus

$$\sum_i \left[\langle \partial \varphi_i | \hat{f} | \varphi_i \rangle - \sum_k \epsilon_{ik} \langle \partial \varphi_i | \varphi_k \rangle \right] = 0$$

$$\langle \partial \varphi_i | \left[\hat{f} | \varphi_i \rangle - \sum_k \epsilon_{ik} | \varphi_k \rangle \right] = 0$$

system of n coupled non-linear equations

$$\hat{f} | \varphi_i \rangle = \sum_k \epsilon_{ik} | \varphi_k \rangle$$

let's try to uncouple them!

Hartree-Fock Theory for n electrons

task: finding the best orbitals that minimize total energy

system of n coupled non-linear equations

$$\hat{f}|\varphi_i\rangle = \sum_k \epsilon_{ik}|\varphi_k\rangle$$

matrix notation

$$\mathbf{F}|\varphi\rangle = \epsilon|\varphi\rangle$$

unitary transformation

$$|\varphi'_a\rangle = \hat{U}|\phi_a\rangle = \sum_k |\varphi_k\rangle U_{ka} \quad \hat{U}^\dagger \hat{U} = \mathbf{1} \quad \hat{U}^\dagger = \hat{U}^{-1}$$

Fock equation

$$\mathbf{U}\mathbf{F}\mathbf{U}^\dagger \mathbf{U}|\varphi\rangle = \mathbf{U}\epsilon\mathbf{U}^\dagger \mathbf{U}|\varphi\rangle$$

choose unitary transformation matrix that diagonalizes ϵ

$$\mathbf{F}'|\varphi'\rangle = \epsilon'|\varphi'\rangle \quad \epsilon'_{ik} = \delta_{ik}\epsilon_k$$

Hartree-Fock Theory for n electrons

task: finding the best orbitals that minimize total energy

canonical Hartree-Fock orbitals

$$\hat{f}|\varphi'_i\rangle = \epsilon'_i|\varphi'_i\rangle$$

orbital energies (ionization energies and electron affinities)

$$\begin{aligned}\epsilon_i = & \langle \varphi_i | \hat{h} | \varphi \rangle + \sum_k \int \frac{|\varphi_i(\mathbf{x}_1)|^2 |\varphi_k(\mathbf{x}_2)|^2}{r_{12}} d\mathbf{x}_1 d\mathbf{x}_2 \\ & - \sum_k \int \frac{\varphi_i^*(\mathbf{x}_1) \varphi_k^*(\mathbf{x}_2) \varphi_k(\mathbf{x}_1) \varphi_i(\mathbf{x}_2)}{r_{12}} d\mathbf{x}_1 d\mathbf{x}_2\end{aligned}$$

$$E = \sum_i h_{ii} + \sum_i \sum_k J_{ik} - \sum_i \sum_k K_{ik}$$

total energy (correct double counting, also first order energy in PT)

$$E = \sum_i h_{ii} + \frac{1}{2} \sum_i \sum_k J_{ik} - \frac{1}{2} \sum_i \sum_k K_{ik}$$

Hartree-Fock Theory for n electrons

task: finding the best orbitals that minimize total energy

canonical Hartree-Fock orbitals

$$\hat{f}|\varphi_i\rangle = \epsilon_i|\varphi_i\rangle$$

orbital energies (ionization energies and electron affinities)

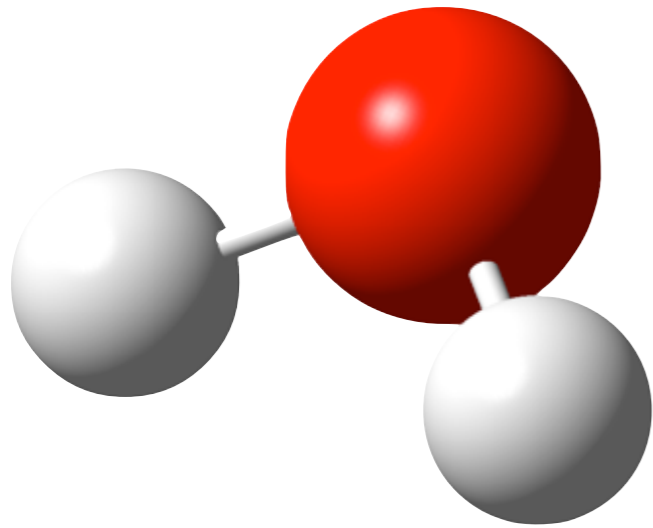
$$\begin{aligned}\epsilon_i = & \langle \varphi_i | \hat{h} | \varphi_i \rangle + \sum_k \int \frac{|\varphi_i(\mathbf{x}_1)|^2 |\varphi_k(\mathbf{x}_2)|^2}{r_{12}} d\mathbf{x}_1 d\mathbf{x}_2 \\ & - \sum_k \int \frac{\varphi_i^*(\mathbf{x}_1) \varphi_k^*(\mathbf{x}_2) \varphi_k(\mathbf{x}_1) \varphi_i(\mathbf{x}_2)}{r_{12}} d\mathbf{x}_1 d\mathbf{x}_2\end{aligned}$$

$$E = \sum_i h_{ii} + \sum_i \sum_k J_{ik} - \sum_i \sum_k K_{ik}$$

total energy (correct double counting)

$$E = \sum_i h_{ii} + \frac{1}{2} \sum_i \sum_k J_{ik} - \frac{1}{2} \sum_i \sum_k K_{ik}$$

Wassermolekül



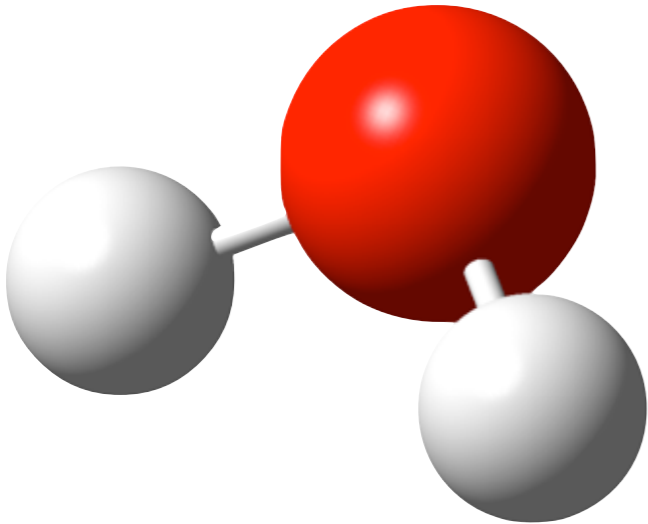
2 H⁺

1 O⁸⁺

10 Elektronen

10 Molekülorbitale

Wassermolekül

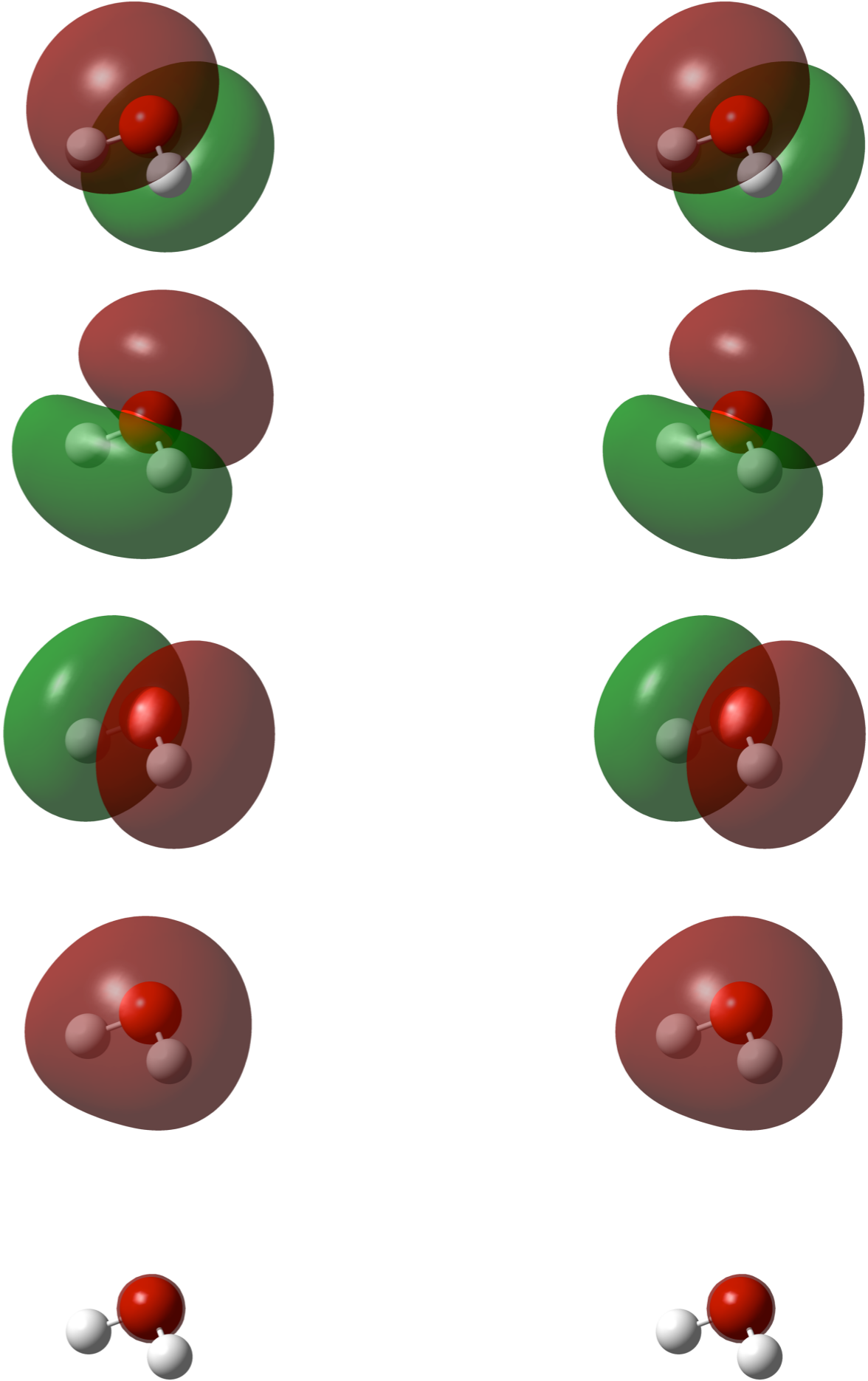


2 H⁺

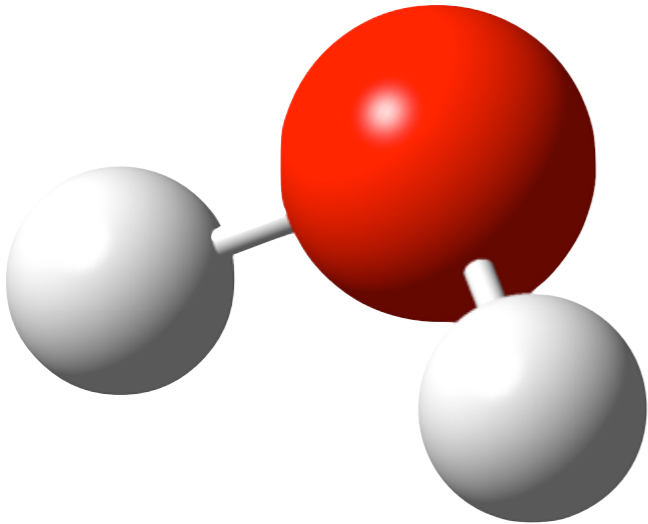
1 O⁸⁺

10 Elektronen

10 Molekülorbitale



Wassermolekül



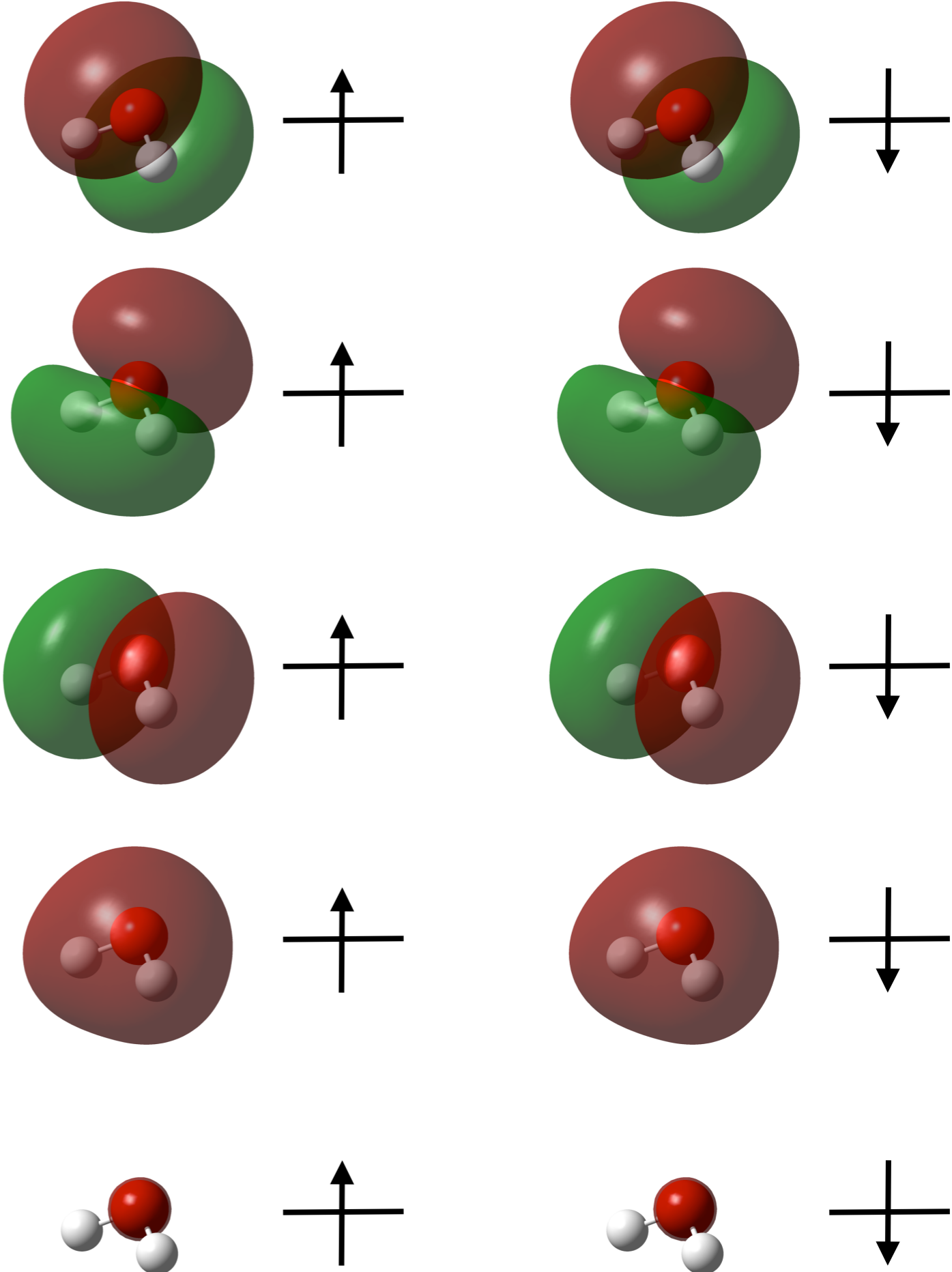
2 H⁺

1 O⁸⁺

10 Elektronen

10 Molekülorbitale

Energie ↑



Hartree-Fock Theory for n electrons

Hartree-Fock eigenvalue equations

$$\hat{f}(\mathbf{r})\varphi_i(\mathbf{x}) = \epsilon_i\varphi(\mathbf{x})$$

solving non-linear eigenvalues equations numerically

step 1: get rid of spin and express in real spatial orbitals

step 2: expand spatial orbitals in basis functions

restricted Hartree Fock

electron pair with opposite spin in same spatial orbital

$$\varphi_i(\mathbf{x}) = \phi_j(\mathbf{r})\alpha(s)$$

$$\varphi_{i+1}(\mathbf{x}) = \phi_j(\mathbf{r})\beta(s)$$

Hartree-Fock Theory for n electrons

solving non-linear eigenvalues equations numerically

step 1: get rid of spin and express in real spatial orbitals

$$\begin{aligned}\hat{f}(\mathbf{x}_1)\phi_i(\mathbf{r}_1)\alpha(s_1) &= \hat{h}^0(\mathbf{r}_1)\phi_i(\mathbf{r}_1)\alpha(s_1) \\ &+ \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2)\alpha^*(s_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2)\alpha(s_2)\phi_i(\mathbf{r}_1)\alpha(s_1) d\mathbf{r}_2 ds_2 \\ &+ \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2)\beta^*(s_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2)\beta(s_2)\phi_i(\mathbf{r}_1)\alpha(s_1) d\mathbf{r}_2 ds_2 \\ &- \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2)\alpha^*(s_2) \frac{1}{r_{12}} \phi_i(\mathbf{r}_2)\alpha(s_2)\phi_k(\mathbf{r}_1)\alpha(s_1) d\mathbf{r}_2 ds_2 \\ &- \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2)\beta^*(s_2) \frac{1}{r_{12}} \phi_i(\mathbf{r}_2)\alpha(s_2)\phi_k(\mathbf{r}_1)\beta(s_1) d\mathbf{r}_2 ds_2\end{aligned}$$

Hartree-Fock Theory for n electrons

solving non-linear eigenvalues equations numerically

step 1: get rid of spin and express in real spatial orbitals

$$\begin{aligned}\hat{f}(\mathbf{x}_1)\phi_i(\mathbf{r}_1)\alpha(s_1) &= \hat{h}^0(\mathbf{r}_1)\phi_i(\mathbf{r}_1)\alpha(s_1) \\ &+ \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2)\alpha^*(s_2)\frac{1}{r_{12}}\phi_k(\mathbf{r}_2)\alpha(s_2)\phi_i(\mathbf{r}_1)\alpha(s_1)d\mathbf{r}_2ds_2 \\ &+ \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2)\beta^*(s_2)\frac{1}{r_{12}}\phi_k(\mathbf{r}_2)\beta(s_2)\phi_i(\mathbf{r}_1)\alpha(s_1)d\mathbf{r}_2ds_2 \\ &- \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2)\alpha^*(s_2)\frac{1}{r_{12}}\phi_i(\mathbf{r}_2)\alpha(s_2)\phi_k(\mathbf{r}_1)\alpha(s_1)d\mathbf{r}_2ds_2 \\ &- \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2)\beta^*(s_2)\frac{1}{r_{12}}\phi_i(\mathbf{r}_2)\alpha(s_2)\phi_k(\mathbf{r}_1)\beta(s_1)d\mathbf{r}_2ds_2\end{aligned}$$

Hartree-Fock Theory for n electrons

solving non-linear eigenvalues equations numerically

step 1: get rid of spin and express in real spatial orbitals

$$\begin{aligned}\hat{f}(\mathbf{x}_1)\phi_i(\mathbf{r}_1)\alpha(s_1) &= \hat{h}^0(\mathbf{r}_1)\phi_i(\mathbf{r}_1)\alpha(s_1) \\ &+ \sum_k^{n/2} \int \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2)\phi_i(\mathbf{r}_1)\alpha(s_1) d\mathbf{r}_2 \\ &+ \sum_k^{n/2} \int \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2)\phi_i(\mathbf{r}_1)\alpha(s_1) d\mathbf{r}_2 \\ &- \sum_k^{n/2} \int \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_i(\mathbf{r}_2)\phi_k(\mathbf{r}_1)\alpha(s_1) d\mathbf{r}_2\end{aligned}$$

Hartree-Fock Theory for n electrons

solving non-linear eigenvalues equations numerically

step 1: get rid of spin and express in real spatial orbitals

$$\begin{aligned} \int \alpha^*(s_1) \hat{f}(\mathbf{x}_1) \alpha(s_1) ds_1 \phi_i(\mathbf{r}_1) &= \int \alpha^*(s_1) \hat{h}^0(\mathbf{r}_1) \phi(\mathbf{r}_1) \alpha(s_1) ds_1 \\ &+ \sum_k^{n/2} \int \int \alpha^*(s_1) \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2) \phi_i(\mathbf{r}_1) \alpha(s_1) d\mathbf{r}_2 ds_1 \\ &+ \sum_k^{n/2} \int \int \alpha^*(s_1) \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2) \phi_i(\mathbf{r}_1) \alpha(s_1) d\mathbf{r}_2 ds_1 \\ &- \sum_k^{n/2} \int \int \alpha^*(s_1) \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_i(\mathbf{r}_2) \phi_k(\mathbf{r}_1) \alpha(s_1) d\mathbf{r}_2 ds_1 \end{aligned}$$

Hartree-Fock eigenvalue equation for spatial orbitals

$$\begin{aligned} \hat{f}(\mathbf{r}_1) \phi_i(\mathbf{r}_1) &= \hat{h}^0(\mathbf{r}_1) \phi(\mathbf{r}_1) \\ &+ 2 \sum_k^{n/2} \int \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2) \phi_i(\mathbf{r}_1) d\mathbf{r}_2 \\ &- \sum_k^{n/2} \int \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_i(\mathbf{r}_2) \phi_k(\mathbf{r}_1) d\mathbf{r} \\ &= \epsilon_i \phi_i(\mathbf{r}_1) \end{aligned}$$

Hartree-Fock Theory for n electrons

solving non-linear eigenvalues equations numerically

step 1: get rid of spin and express in real spatial orbitals

$$\begin{aligned}\hat{f}(\mathbf{r}_1)\phi_i(\mathbf{r}_1) &= \hat{h}^0(\mathbf{r}_1)\phi(\mathbf{r}_1) \\ &+ 2 \sum_k^{n/2} \int \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2) \phi_i(\mathbf{r}_1) d\mathbf{r}_2 \\ &- \sum_k^{n/2} \int \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_i(\mathbf{r}_2) \phi_k(\mathbf{r}_1) d\mathbf{r} \\ &= \epsilon_i \phi_i(\mathbf{r}_1)\end{aligned}$$

step 2: expand spatial orbitals in basis functions (basisset)

$$\phi_i(\mathbf{r}) = \sum_j^k c_{ij} \gamma_j(\mathbf{r} - \mathbf{R}_j)$$

Hartree-Fock Theory for n electrons

linear combination of atomic orbitals

$$\phi_i(\mathbf{r}) = \sum_j^k c_{ij} \gamma_j(\mathbf{r} - \mathbf{R}_j)$$

hydrogen-like orbitals (one possibility out of many...)

$$\gamma_1 = \psi^{1s}(\zeta_1)$$

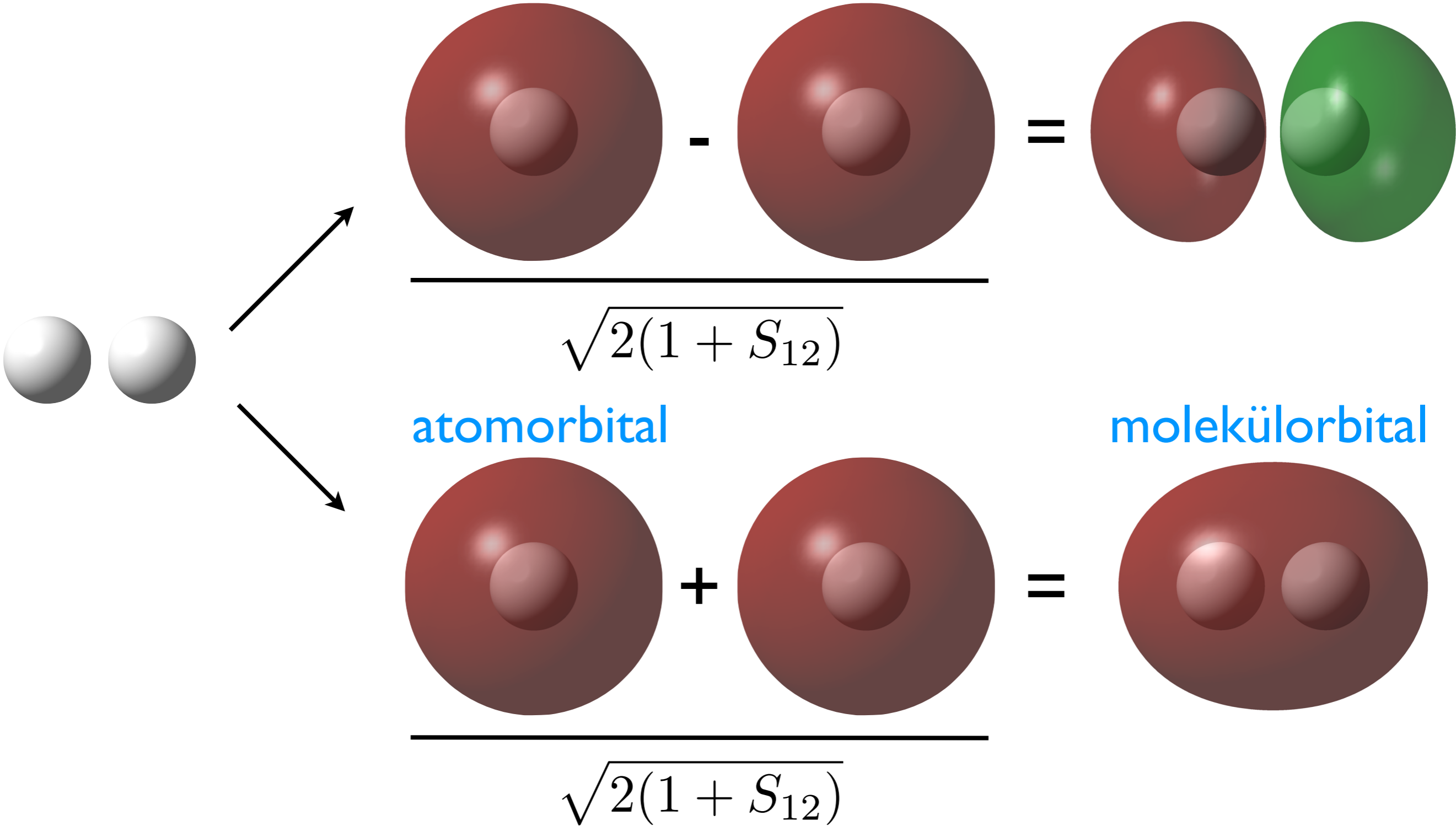
$$\gamma_2 = \psi^{2s}(\zeta_2)$$

$$\gamma_3 = \psi^{2p}(\zeta_3)$$

$$\gamma_4 = \dots$$

Wasserstoffmolekül

Lineare Kombination von einzelne Wasserstoff-Orbitale



Hartree-Fock Theory for n electrons

solving non-linear eigenvalues equations numerically

step 2: expand spatial orbitals in basis functions

$$\hat{f}(\mathbf{r}_1)\phi_i(\mathbf{r}_1) = \epsilon_i\phi_i(\mathbf{r}_1)$$

$$\phi_i(\mathbf{r}) = \sum_j^k c_{ij}\gamma_j(\mathbf{r} - \mathbf{R}_j)$$

$$\hat{f}(\mathbf{r}_1) \sum_{\nu} c_{\nu i}\gamma_{\nu}(\mathbf{r}_1) = \epsilon_i \sum_{\nu} c_{\nu i}\gamma_{\nu}(\mathbf{r}_1)$$

$$\sum_{\nu} c_{\nu i} \int \gamma_{\mu}^*(\mathbf{r}_1)\hat{f}(\mathbf{r}_1)\gamma_{\nu}(\mathbf{r}_1)d\mathbf{r}_1 = \epsilon_i \sum_{\nu} c_{\nu i} \int \gamma_{\mu}(\mathbf{r}_1)\gamma_{\nu}(\mathbf{r}_1)d\mathbf{r}_1$$

Hartree-Fock Theory for n electrons

solving non-linear eigenvalues equations numerically

step 2: expand spatial orbitals in basis functions (basisset)

$$\sum_{\nu} c_{\nu i} \int \gamma_{\mu}^*(\mathbf{r}_1) \hat{f}(\mathbf{r}_1) \gamma_{\nu}(\mathbf{r}_1) d\mathbf{r}_1 = \epsilon_i \sum_{\nu} c_{\nu i} \int \gamma_{\mu}(\mathbf{r}_1) \gamma_{\nu}(\mathbf{r}_1) d\mathbf{r}_1$$

express in terms of matrices

$$\sum_{\nu} F_{\mu\nu} c_{\nu i} = \epsilon_i \sum_{\nu} S_{\mu\nu} c_{\nu i}$$

$$\mathbf{FC} = \mathbf{SC}\epsilon$$

solution if, and only if

$$|\mathbf{F} - \epsilon_i \mathbf{S}| = 0$$

Hartree-Fock Theory for n electrons

solving non-linear eigenvalues equations numerically

non-linear: \mathbf{F} depends on \mathbf{C}

$$\begin{aligned} F_{\mu\nu} &= \int \gamma_{\mu}^*(\mathbf{r}_1) \hat{h}^0(\mathbf{r}_1) \gamma_{\nu}(\mathbf{r}_1) d\mathbf{r}_1 \\ &+ 2 \sum_a \int \int \gamma_{\mu}^*(\mathbf{r}_1) \phi_a^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_a(\mathbf{r}_2) \gamma_{\nu}(\mathbf{r}_1) d\mathbf{r}_1 d\mathbf{r} \\ &- \sum_a \int \int \gamma_{\mu}^*(\mathbf{r}_1) \phi_a^*(\mathbf{r}_2) \frac{1}{r_{12}} \gamma_{\nu}(\mathbf{r}_2) \phi_a(\mathbf{r}_1) d\mathbf{r}_1 d\mathbf{r} \end{aligned}$$

$$\begin{aligned} F_{\mu\nu} &= h_{\mu\nu}^0 \\ &+ 2 \sum_a \sum_{\kappa} \sum_{\lambda} c_{\kappa a}^* c_{\lambda a}^* \int \int \gamma_{\mu}^*(\mathbf{r}_1) \gamma_{\kappa}^*(\mathbf{r}_2) \frac{1}{r_{12}} \gamma_{\lambda}(\mathbf{r}_2) \gamma_{\nu}(\mathbf{r}_1) d\mathbf{r}_1 d\mathbf{r} \\ &- \sum_a \sum_{\kappa} \sum_{\lambda} c_{\kappa a}^* c_{\lambda a}^* \int \int \gamma_{\mu}^*(\mathbf{r}_1) \gamma_{\kappa}^*(\mathbf{r}_2) \frac{1}{r_{12}} \gamma_{\nu}(\mathbf{r}_2) \gamma_a(\mathbf{r}_1) d\mathbf{r}_1 d\mathbf{r} \end{aligned}$$

Hartree-Fock

Roothaan-Hall equations

non-linear eigenvalue problem

$$\mathbf{F}\mathbf{c} = \epsilon_i\mathbf{S}\mathbf{c}$$

practical algorithm

iterate until self-consistency

pre-compute integrals of basisset

$$S_{\mu\nu} = \langle \gamma_\mu | \gamma_\nu \rangle$$

