

# Hartree-Fock Theory for $n$ electrons

Expectations values for one and two electron operators

Hartree product (no spin)

$$\langle \hat{O}_1 \rangle = \sum_a \int \phi_a^*(\mathbf{x}_1) \hat{o}(\mathbf{r}_1) \phi_a(\mathbf{r}_1) d\mathbf{r}_1$$

$$\langle \hat{O}_2 \rangle = \frac{1}{2} \sum_a \sum_b \int \int \phi_a^*(\mathbf{r}_1) \phi_b^*(\mathbf{r}_2) \hat{o}(\mathbf{r}_1, \mathbf{r}_2) \phi_a(\mathbf{r}_1) \phi_b(\mathbf{r}_2) d\mathbf{r}_1 d\mathbf{r}_2$$

Slater determinant (spin, Pauli principle)

$$\langle \hat{O}_1 \rangle = \sum_a \int \phi_a^*(\mathbf{x}_1) \hat{o}(\mathbf{r}_1) \phi_a(\mathbf{x}_1) d\mathbf{x}_1$$

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$$- \frac{1}{2} \sum_a \sum_b \int \int \phi_a^*(\mathbf{x}_1) \phi_b^*(\mathbf{x}_2) \hat{o}(\mathbf{r}_1, \mathbf{r}_2) \phi_b(\mathbf{x}_1) \phi_a(\mathbf{x}_2) d\mathbf{x}_1 d\mathbf{x}_2$$

$$\langle \hat{O}_2 \rangle = \frac{1}{2} \sum_a \sum_b \int \int \phi_a^*(\mathbf{x}_1) \phi_b^*(\mathbf{x}_2) \hat{o}(\mathbf{r}_1, \mathbf{r}_2) (1 - \hat{p}_{12}) \phi_a(\mathbf{x}_1) \phi_b(\mathbf{x}_2) d\mathbf{x}_1 d\mathbf{x}_2$$

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# Hartree-Fock Theory for $n$ electrons

task: finding the best spinorbitals that minimize total energy

Variational principle

$$E \geq E^0$$

minimum

$$\partial E = 0$$

when for all spin orbitals

$$\varphi_i(\mathbf{x}) = \varphi_i(\mathbf{x}) + \partial\varphi_i(\mathbf{x})$$

constraint

$$\langle \varphi_i(\mathbf{x}) | \varphi_k(\mathbf{x}) \rangle = \delta_{ik}$$

Langrange multipliers

# Hartree-Fock Theory for $n$ electrons

task: finding the best orbitals that minimize total energy

Lagrangian (adding a smart zero, or actually  $n^2$  smart zeros)

$$\mathcal{L} = E - \sum_i \sum_k \epsilon_{ki} [\langle \varphi_i | \varphi_k \rangle - \delta_{ik}]$$

minimum

$$\partial \mathcal{L} = \partial E - \sum_i \sum_k \epsilon_{ki} [\langle \partial \varphi_i | \varphi_k \rangle + \langle \varphi_i | \partial \varphi_k \rangle] = 0$$

Lagrangian and energy are real

$$\mathcal{L} = \mathcal{L}^* \quad E = E^*$$

therefore

$$(\sum_i \sum_k \epsilon_{ki} [\langle \varphi_i | \varphi_k \rangle - \delta_{ik}])^* = \sum_i \sum_k \epsilon_{ki} [\langle \varphi_i | \varphi_k \rangle - \delta_{ik}]$$

$$\sum_i \sum_k \epsilon_{ki}^* [\langle \varphi_i | \varphi_k \rangle^* - \delta_{ik}] = \sum_i \sum_k \epsilon_{ki} [\langle \varphi_i | \varphi_k \rangle - \delta_{ik}]$$

$$\sum_i \sum_k \epsilon_{ik}^* [\langle \varphi_k | \varphi_i \rangle - \delta_{ik}] = \sum_i \sum_k \epsilon_{ki} [\langle \varphi_i | \varphi_k \rangle - \delta_{ik}]$$

$$\sum_k \sum_i \epsilon_{ki}^* [\langle \varphi_i | \varphi_k \rangle - \delta_{ik}] = \sum_i \sum_k \epsilon_{ki} [\langle \varphi_i | \varphi_k \rangle - \delta_{ik}]$$

$$\epsilon_{ki}^* = \epsilon_{ki}$$

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$$\partial \mathcal{L} = \partial E - \sum_i \sum_k \epsilon_{ki} [\langle \partial \varphi_i | \varphi_k \rangle + \langle \varphi_i | \partial \varphi_k \rangle] = 0$$

rearranging

$$\begin{aligned} \partial E &= \sum_i \langle \partial \varphi_i | \hat{h} | \varphi_i \rangle + \frac{1}{2} \sum_i \sum_k & & [\langle \partial \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_i | \varphi_k \rangle + \langle \varphi_i | \partial \varphi_k | \frac{1}{r_{12}} | \varphi_i | \varphi_k \rangle \\ &&& - \langle \partial \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_k | \varphi_i \rangle - \langle \varphi_i | \partial \varphi_k | \frac{1}{r_{12}} | \varphi_k | \varphi_i \rangle] \\ &+ \sum_i \langle \varphi_i | \hat{h} | \partial \varphi_i \rangle + \frac{1}{2} \sum_i \sum_k & & [\langle \varphi_i | \varphi_k | \frac{1}{r_{12}} | \partial \varphi_i | \varphi_k \rangle + \langle \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_i | \partial \varphi_k \rangle \\ &&& - \langle \varphi_i | \varphi_k | \frac{1}{r_{12}} | \partial \varphi_k | \varphi_i \rangle - \langle \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_k | \partial \varphi_i \rangle] \end{aligned}$$

# Hartree-Fock Theory for $n$ electrons

task: finding the best orbitals that minimize total energy

Lagrangian (adding a smart zero!)

$$\mathcal{L} = E - \sum_i \sum_k \epsilon_{ki} [\langle \varphi_i | \varphi_k \rangle - \delta_{ik}]$$

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rearranging

$$\begin{aligned} \partial E = & \sum_i \langle \partial \varphi_i | \hat{h} | \varphi_i \rangle + \frac{1}{2} \sum_i \sum_k & [\langle \partial \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_i | \varphi_k \rangle + \langle \varphi_i | \partial \varphi_k | \frac{1}{r_{12}} | \varphi_i | \varphi_k \rangle \\ & & - \langle \partial \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_k | \varphi_i \rangle - \langle \varphi_i | \partial \varphi_k | \frac{1}{r_{12}} | \varphi_k | \varphi_i \rangle] \end{aligned}$$

$$\begin{aligned} & + \sum_i \langle \varphi_i | \hat{h} | \partial \varphi_i \rangle + \frac{1}{2} \sum_i \sum_k & [\langle \varphi_i | \varphi_k | \frac{1}{r_{12}} | \partial \varphi_i | \varphi_k \rangle + \langle \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_i | \partial \varphi_k \rangle \\ & & - \langle \varphi_i | \varphi_k | \frac{1}{r_{12}} | \partial \varphi_k | \varphi_i \rangle - \langle \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_k | \partial \varphi_i \rangle] \end{aligned}$$

complex conjugate

# Hartree-Fock Theory for $n$ electrons

task: finding the best orbitals that minimize total energy

Lagrangian (adding a smart zero!)

$$\mathcal{L} = E - \sum_i \sum_k \epsilon_{ki} [\langle \varphi_i | \varphi_k \rangle - \delta_{ik}]$$

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$$\partial \mathcal{L} = \partial E - \sum_i \sum_k \epsilon_{ki} [\langle \partial \varphi_i | \varphi_k \rangle + \langle \varphi_i | \partial \varphi_k \rangle] = 0$$

rearranging

$$\partial E = \sum_i \langle \partial \varphi_i | \hat{h} | \varphi_i \rangle + \frac{1}{2} \sum_i \sum_k$$

changing variables

$$[\langle \partial \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_i \varphi_k \rangle + \langle \varphi_i | \partial \varphi_k | \frac{1}{r_{12}} | \varphi_i \varphi_k \rangle \\ - \langle \partial \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_k \varphi_i \rangle - \langle \varphi_i | \partial \varphi_k | \frac{1}{r_{12}} | \varphi_k \varphi_i \rangle]$$

$$+ \sum_i \langle \varphi_i | \hat{h} | \partial \varphi_i \rangle + \frac{1}{2} \sum_i \sum_k [\langle \varphi_i \varphi_k | \frac{1}{r_{12}} | \partial \varphi_i \varphi_k \rangle + \langle \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_i \partial \varphi_k \rangle]$$

$$- \langle \varphi_i \varphi_k | \frac{1}{r_{12}} | \partial \varphi_k \varphi_i \rangle - \langle \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_k \partial \varphi_i \rangle]$$

complex conjugate

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Lagrangian (adding a smart zero!)

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$$\partial \mathcal{L} = \partial E - \sum_i \sum_k \epsilon_{ki} [\langle \partial \varphi_i | \varphi_k \rangle + \langle \varphi_i | \partial \varphi_k \rangle] = 0$$

rearranging

$$\begin{aligned} \partial E &= \sum_i \langle \partial \varphi_i | \hat{h} | \varphi_i \rangle + \sum_i \sum_k [\langle \partial \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_i | \varphi_k \rangle - \langle \partial \varphi_i | \varphi_k | \frac{1}{r_{12}} | \varphi_k | \varphi_i \rangle] \\ &\quad + \text{c.c} \end{aligned}$$

# Hartree-Fock Theory for $n$ electrons

task: finding the best orbitals that minimize total energy

rearranging

$$\begin{aligned}\partial E = & \sum_i \langle \partial \varphi_i | \hat{h} | \varphi_i \rangle + \sum_i \sum_k \langle \partial \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_i \varphi_k \rangle \\ & - \sum_i \sum_k \langle \partial \varphi_i \varphi_k | \frac{1}{r_{12}} | \varphi_k \varphi_i \rangle \\ & + \text{c.c.}\end{aligned}$$

coulomb

exchange

for each orbital:

$$\langle \partial \varphi_i | \hat{h} | \varphi_i \rangle + \langle \partial \varphi_i | \sum_k \frac{|\varphi_k|^2}{r_{12}} | \varphi_i \rangle - \langle \partial \varphi_i | \sum_k \langle \varphi_k | \frac{1}{r_{12}} | \varphi_i \rangle \varphi_k \rangle$$

# Hartree-Fock Theory for $n$ electrons

task: finding the best orbitals that minimize total energy

rearranging a bit more

$$\langle \partial\varphi_i | \hat{h} | \varphi_i \rangle + \langle \partial\varphi_i | \sum_k \frac{|\varphi_k|^2}{r_{12}} | \varphi_i \rangle - \langle \partial\varphi_i | \sum_k \langle \varphi_k | \frac{1}{r_{12}} | \varphi_i \rangle \varphi_k \rangle$$

local coulomb operator

$$\hat{J}_k |\varphi_i(1)\rangle = \frac{|\varphi_k|^2}{r_{12}} |\varphi_i(1)\rangle$$

non-local exchange operator

$$\hat{K}_k |\varphi_i(1)\rangle = \langle \varphi_k(2) | \frac{1}{r_{12}} | \varphi_i(2) \rangle | \varphi_k(1) \rangle$$

Fock operator

$$\hat{f} = \hat{h} + \sum_k [\hat{J}_k - \hat{K}_k]$$

# Hartree-Fock Theory for $n$ electrons

task: finding the best orbitals that minimize total energy

going back the the lagrangian

$$\partial \mathcal{L} = \partial E - \sum_i \sum_k \epsilon_{ik} [\langle \partial \varphi_i | \varphi_k \rangle + \langle \varphi_i | \partial \varphi_k \rangle] = 0$$

thus

$$\sum_i \left[ \langle \partial \varphi_i | \hat{f} | \varphi_i \rangle - \sum_k \epsilon_{ki} \langle \partial \varphi_i | \varphi_k \rangle \right] + c.c. = 0$$

$$\langle \partial \varphi_i | \left[ \hat{f} | \varphi_i \rangle - \sum_k \epsilon_{ki} \varphi_k \rangle \right] + c.c. = 0$$

system of  $n$  coupled non-linear equations

$$\hat{f} | \varphi_i \rangle = \sum_k \epsilon_{ki} | \varphi_k \rangle$$

let's try to uncouple them!

# Hartree-Fock Theory for $n$ electrons

task: finding the best orbitals that minimize total energy

system of  $n$  coupled non-linear equations

$$\hat{f}|\varphi_i\rangle = \sum_k \epsilon_{ki} |\varphi_k\rangle$$

matrix notation

$$\mathbf{F}|\varphi\rangle = \boldsymbol{\epsilon}|\varphi\rangle$$

unitary transformation

$$|\varphi_\alpha\rangle = \sum_i |\varphi_i\rangle \langle \varphi_i| \varphi_\alpha\rangle = \sum_i |\varphi_i\rangle U_{i\alpha} \quad \hat{U}^\dagger \hat{U} = 1 \quad \hat{U}^\dagger = \hat{U}^{-1}$$

Fock equation

$$\mathbf{U}\mathbf{F}\mathbf{U}^\dagger \mathbf{U}|\varphi\rangle = \mathbf{U}\boldsymbol{\epsilon}\mathbf{U}^\dagger \mathbf{U}|\varphi\rangle$$

choose unitary transformation matrix that diagonalizes  $\boldsymbol{\epsilon}$

$$\mathbf{F}'|\varphi'\rangle = \boldsymbol{\epsilon}'|\varphi'\rangle \quad \epsilon'_{ik} = \delta_{ik} \epsilon_k$$

# Hartree-Fock Theory for $n$ electrons

task: finding the best orbitals that minimize total energy

canonical Hartree-Fock orbitals

$$\hat{f}|\varphi'_i\rangle = \epsilon'_i |\varphi'_i\rangle$$

orbital energies (ionization energies and electron affinities)

$$\epsilon_i = \langle \varphi_i | \hat{h} | \varphi \rangle + \sum_k \int \frac{|\varphi_i(\mathbf{x})_1|^2 |\varphi_k(\mathbf{x}_2)|^2}{r_{12}} d\mathbf{x}_1 d\mathbf{x}_2$$

$$- \sum_k \int \frac{\varphi_i^*(\mathbf{x}_1) \varphi_k^*(\mathbf{x}_2) \varphi_k(\mathbf{x}_1) \varphi_i(\mathbf{x}_2)}{r_{12}} d\mathbf{x}_1 d\mathbf{x}_2$$

$$E = \sum_i h_{ii} + \sum_i \sum_k J_{ik} - \sum_i \sum_k K_{ik}$$

total energy (correct double counting, also first order energy in PT)

$$E = \sum_i h_{ii} + \frac{1}{2} \sum_i \sum_k J_{ik} - \frac{1}{2} \sum_i \sum_k K_{ik}$$

# Hartree-Fock Theory for $n$ electrons

task: finding the best orbitals that minimize total energy

canonical Hartree-Fock orbitals

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$$\epsilon_i = \langle \varphi_i | \hat{h} | \varphi \rangle + \sum_k \int \frac{|\varphi_i(\mathbf{x})_1|^2 |\varphi_k(\mathbf{x}_2)|^2}{r_{12}} d\mathbf{x}_1 d\mathbf{x}_2$$

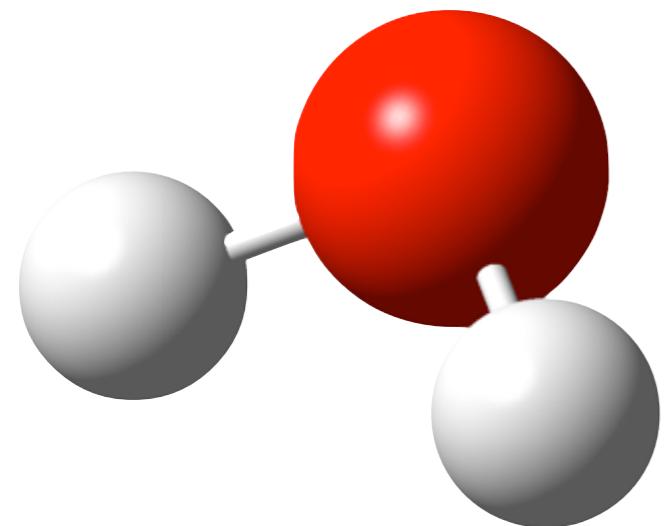
$$- \sum_k \int \frac{\varphi_i^*(\mathbf{x}_1) \varphi_k^*(\mathbf{x}_2) \varphi_k(\mathbf{x}_1) \varphi_i(\mathbf{x}_2)}{r_{12}} d\mathbf{x}_1 d\mathbf{x}_2$$

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# Wassermolekül



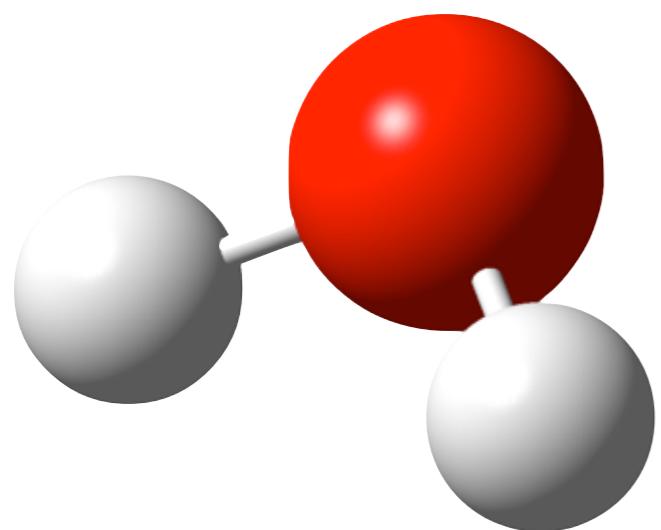
2 H<sup>+</sup>

1 O<sup>8+</sup>

10 Elektronen

10 Molekülorbitale

# Wassermolekül

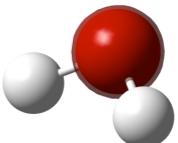
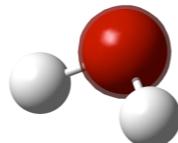
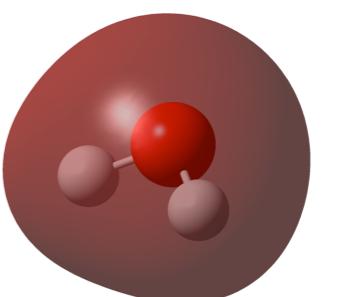
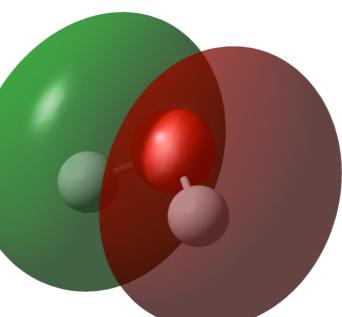
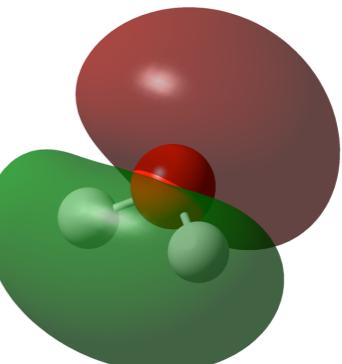
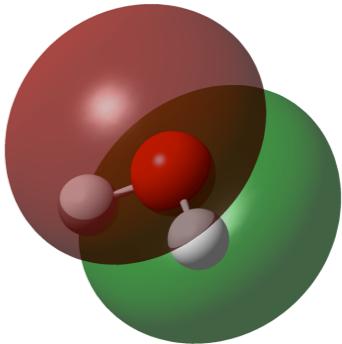
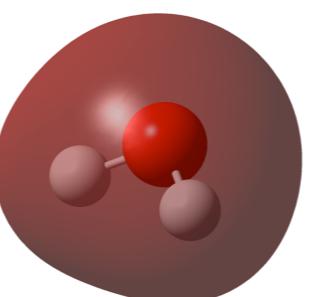
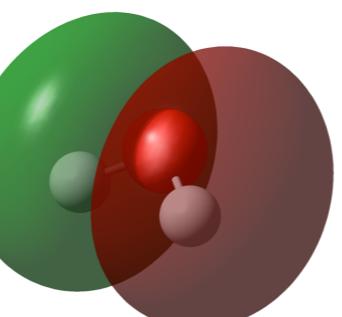
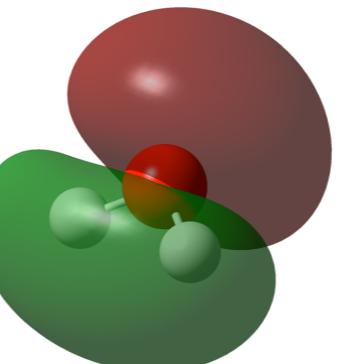
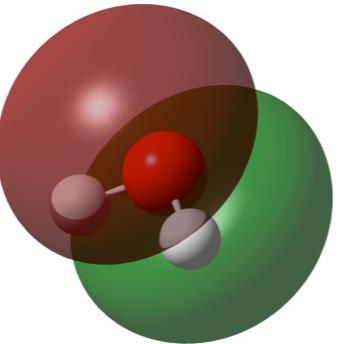


2 H<sup>+</sup>

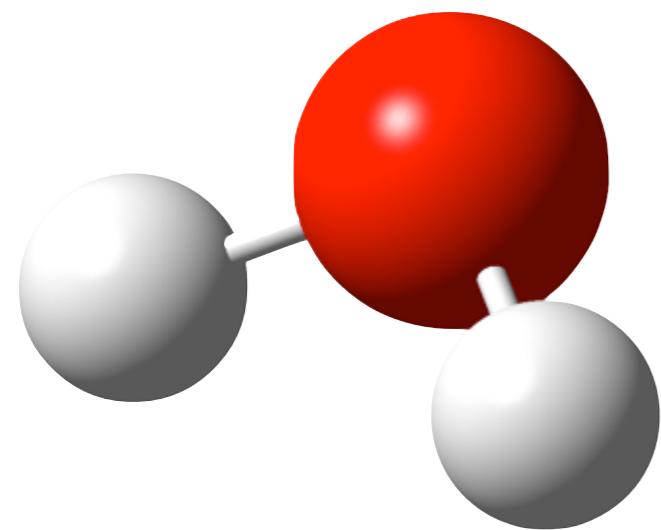
1 O<sup>8+</sup>

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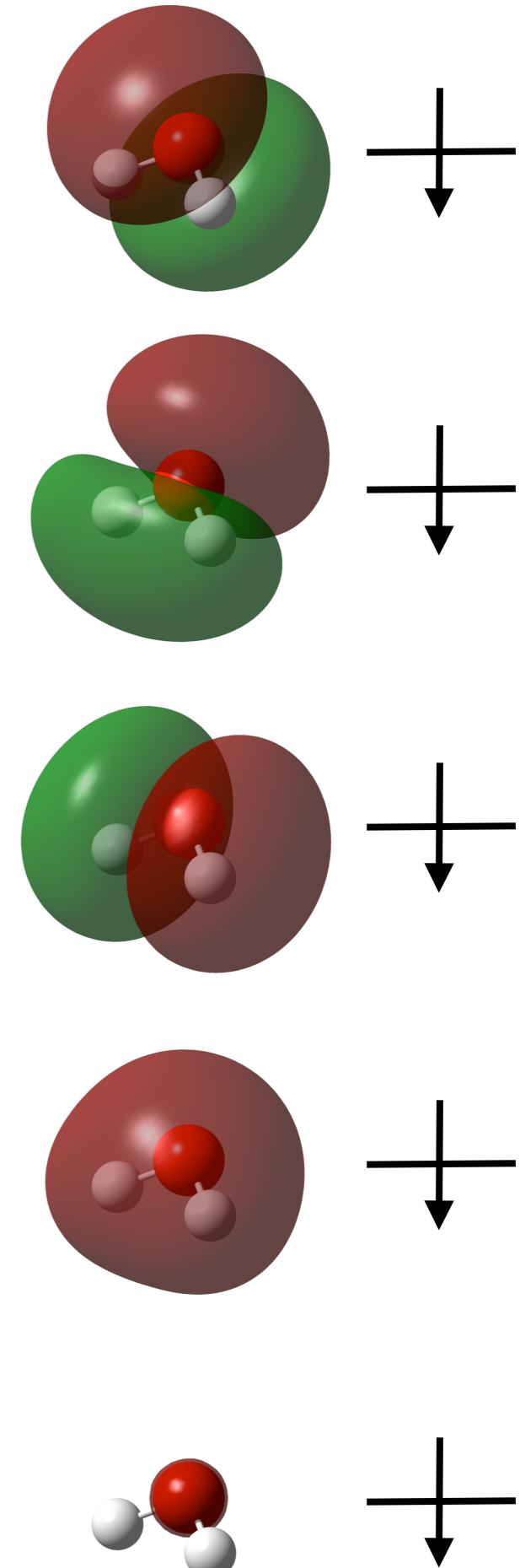
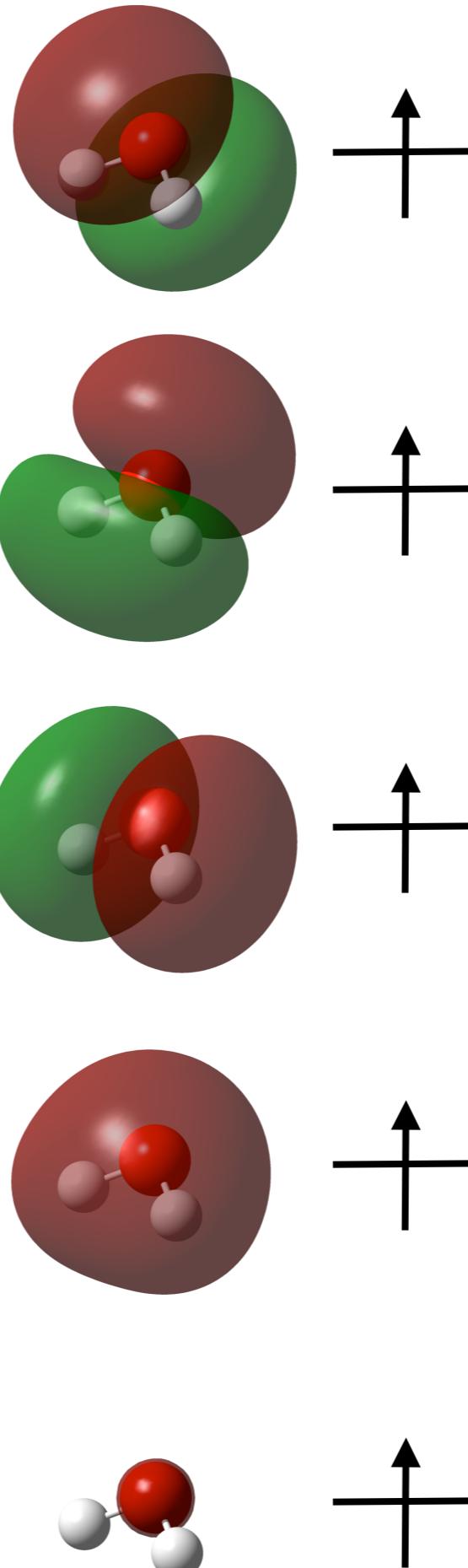


2 H<sup>+</sup>

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# Hartree-Fock Theory for $n$ electrons

Hartree-Fock eigenvalue equations

$$\hat{f}(\mathbf{r})\varphi_i(\mathbf{x}) = \epsilon_i\varphi(\mathbf{x})$$

solving non-linear eigenvalues equations numerically

step 1: get rid of spin and express in real spatial orbitals

step 2: expand spatial orbitals in basis functions

restricted Hartree Fock

electron pair with opposite spin in same spatial orbital

$$\varphi_i(\mathbf{x}) = \phi_j(\mathbf{r})\alpha(s)$$

$$\varphi_{i+1}(\mathbf{x}) = \phi_j(\mathbf{r})\beta(s)$$

# Hartree-Fock Theory for $n$ electrons

solving non-linear eigenvalues equations numerically

step I: get rid of spin and express in real spatial orbitals

$$\begin{aligned}\hat{f}(\mathbf{x}_1)\phi_i(\mathbf{r}_1)\alpha(s_1) &= \hat{h}^0(\mathbf{r}_1)\phi_i(\mathbf{r}_1)\alpha(s_1) \\ &\quad + \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2)\alpha^*(s_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2)\alpha(s_2)\phi_i(\mathbf{r}_1)\alpha(s_1) d\mathbf{r}_2 ds_2 \\ &\quad + \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2)\beta^*(s_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2)\beta(s_2)\phi_i(\mathbf{r}_1)\alpha(s_1) d\mathbf{r}_2 ds_2 \\ &\quad - \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2)\alpha^*(s_2) \frac{1}{r_{12}} \phi_i(\mathbf{r}_2)\alpha(s_2)\phi_k(\mathbf{r}_1)\alpha(s_1) d\mathbf{r}_2 ds_2 \\ &\quad - \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2)\beta^*(s_2) \frac{1}{r_{12}} \phi_i(\mathbf{r}_2)\alpha(s_2)\phi_k(\mathbf{r}_1)\beta(s_1) d\mathbf{r}_2 ds_2\end{aligned}$$

# Hartree-Fock Theory for $n$ electrons

solving non-linear eigenvalues equations numerically

step I: get rid of spin and express in real spatial orbitals

$$\hat{f}(\mathbf{x}_1)\phi_i(\mathbf{r}_1)\alpha(s_1) = \hat{h}^0(\mathbf{r}_1)\phi_i(\mathbf{r}_1)\alpha(s_1)$$

$$+ \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2) \alpha^*(s_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2) \alpha(s_2) \phi_i(\mathbf{r}_1) \alpha(s_1) d\mathbf{r}_2 ds_2$$

$$+ \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2) \beta^*(s_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2) \beta(s_2) \phi_i(\mathbf{r}_1) \alpha(s_1) d\mathbf{r}_2 ds_2$$

$$- \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2) \alpha^*(s_2) \frac{1}{r_{12}} \phi_i(\mathbf{r}_2) \alpha(s_2) \phi_k(\mathbf{r}_1) \alpha(s_1) d\mathbf{r}_2 ds_2$$

$$- \sum_k^{n/2} \int \int \phi_k^*(\mathbf{r}_2) \beta^*(s_2) \frac{1}{r_{12}} \phi_i(\mathbf{r}_2) \alpha(s_2) \phi_k(\mathbf{r}_1) \beta(s_1) d\mathbf{r}_2 ds_2$$

# Hartree-Fock Theory for $n$ electrons

solving non-linear eigenvalues equations numerically

step I: get rid of spin and express in real spatial orbitals

$$\hat{f}(\mathbf{x}_1)\phi_i(\mathbf{r}_1)\alpha(s_1) = \hat{h}^0(\mathbf{r}_1)\phi_i(\mathbf{r}_1)\alpha(s_1)$$

$$+ \sum_k^{n/2} \int \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2) \phi_i(\mathbf{r}_1) \alpha(s_1) d\mathbf{r}_2$$

$$+ \sum_k^{n/2} \int \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2) \phi_i(\mathbf{r}_1) \alpha(s_1) d\mathbf{r}_2$$

$$- \sum_k^{n/2} \int \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_i(\mathbf{r}_2) \phi_k(\mathbf{r}_1) \alpha(s_1) d\mathbf{r}_2$$

# Hartree-Fock Theory for $n$ electrons

solving non-linear eigenvalues equations numerically

step I: get rid of spin and express in real spatial orbitals

$$\begin{aligned} \int \alpha^*(s_1) \hat{f}(\mathbf{x}_1) \alpha(s_1) ds_1 \phi_i(\mathbf{r}_1) &= \int \alpha^*(s_1) \hat{h}^0(\mathbf{r}_1) \phi(\mathbf{r}_1) \alpha(s_1) ds_1 \\ &\quad + \sum_k^{n/2} \int \int \alpha^*(s_1) \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2) \phi_i(\mathbf{r}_1) \alpha(s_1) d\mathbf{r}_2 ds_1 \\ &\quad + \sum_k^{n/2} \int \int \alpha^*(s_1) \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2) \phi_i(\mathbf{r}_1) \alpha(s_1) d\mathbf{r}_2 ds_1 \\ &\quad - \sum_k^{n/2} \int \int \alpha^*(s_1) \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_i(\mathbf{r}_2) \phi_k(\mathbf{r}_1) \alpha(s_1) d\mathbf{r}_2 ds_1 \end{aligned}$$

Hartree-Fock eigenvalue equation for spatial orbitals

$$\begin{aligned} \hat{f}(\mathbf{r}_1) \phi_i(\mathbf{r}_1) &= \hat{h}^0(\mathbf{r}_1) \phi(\mathbf{r}_1) \\ &\quad + 2 \sum_k^{n/2} \int \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2) \phi_i(\mathbf{r}_1) d\mathbf{r}_2 \\ &\quad - \sum_k^{n/2} \int \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_i(\mathbf{r}_2) \phi_k(\mathbf{r}_1) d\mathbf{r}_2 \\ &= \epsilon_i \phi_i(\mathbf{r}_1) \end{aligned}$$

# Hartree-Fock Theory for $n$ electrons

solving non-linear eigenvalues equations numerically

step 1: get rid of spin and express in real spatial orbitals

$$\begin{aligned}\hat{f}(\mathbf{r}_1)\phi_i(\mathbf{r}_1) &= \hat{h}^0(\mathbf{r}_1)\phi_i(\mathbf{r}_1) \\ &\quad + 2 \sum_k^{n/2} \int \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_k(\mathbf{r}_2) \phi_i(\mathbf{r}_1) d\mathbf{r}_2 \\ &\quad - \sum_k^{n/2} \int \phi_k^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_i(\mathbf{r}_2) \phi_k(\mathbf{r}_1) d\mathbf{r}_2 \\ &= \epsilon_i \phi_i(\mathbf{r}_1)\end{aligned}$$

step 2: expand spatial orbitals in basis functions (basisset)

$$\phi_i(\mathbf{r}) = \sum_j^k c_{ij} \gamma_j(\mathbf{r} - \mathbf{R}_j)$$

# Hartree-Fock Theory for $n$ electrons

linear combination of atomic orbitals

$$\phi_i(\mathbf{r}) = \sum_j^k c_{ij} \gamma_j(\mathbf{r} - \mathbf{R}_j)$$

hydrogen-like orbitals (one possibility out of many....)

$$\gamma_1 = \psi^{1s}(\zeta_1)$$

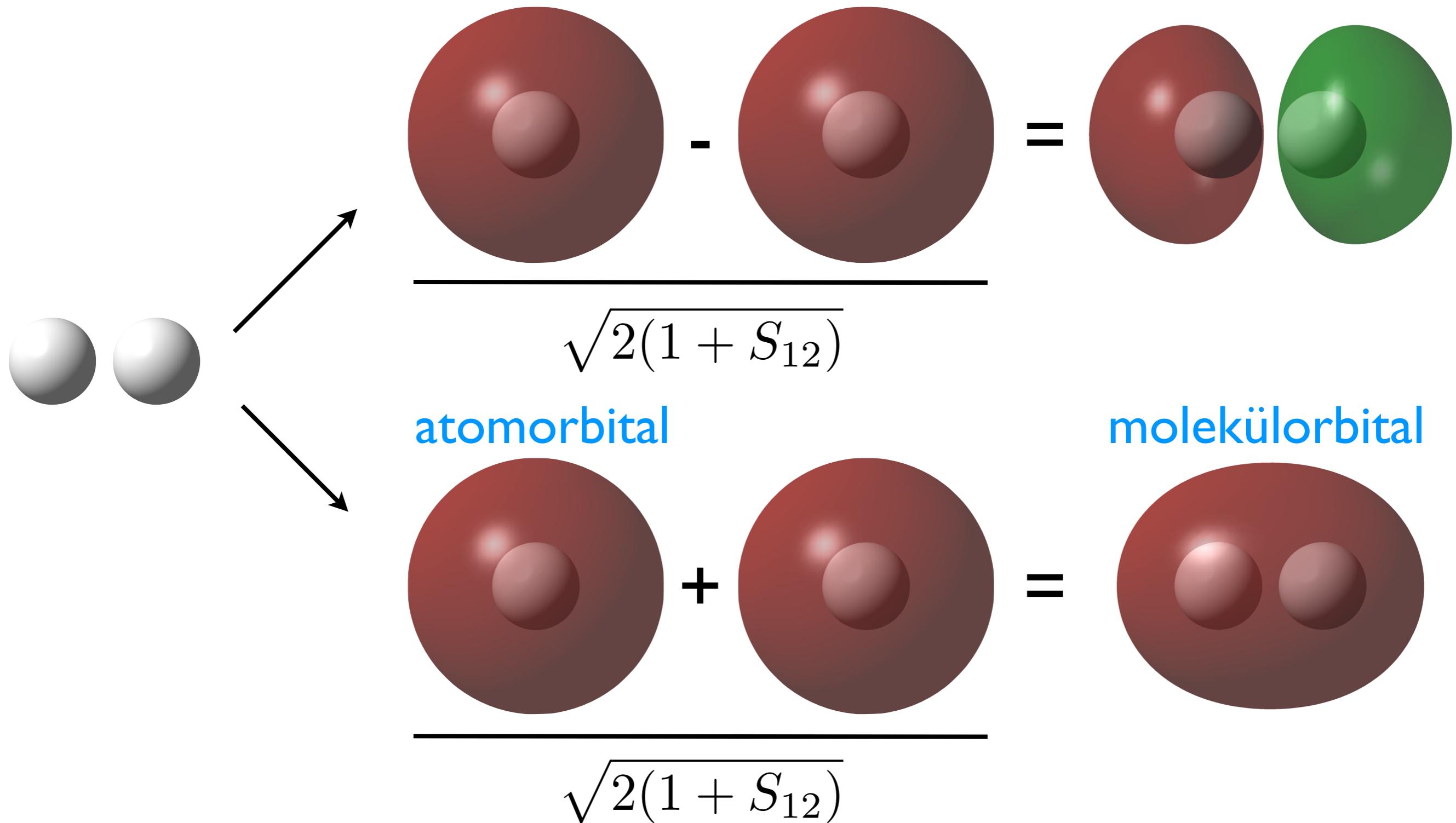
$$\gamma_2 = \psi^{2s}(\zeta_2)$$

$$\gamma_3 = \psi^{2p}(\zeta_3)$$

$$\gamma_4 = \dots$$

# Wasserstoffmolekül

Lineare Kombination von einzelne Wasserstoff-Orbitale



# Hartree-Fock Theory for $n$ electrons

solving non-linear eigenvalues equations numerically

step 2: expand spatial orbitals in basis functions

$$\hat{f}(\mathbf{r}_1)\phi_i(\mathbf{r}_1) = \epsilon_i\phi_i(\mathbf{r}_1)$$

$$\phi_i(\mathbf{r}) = \sum_j^k c_{ij}\gamma_j(\mathbf{r} - \mathbf{R}_j)$$

$$\hat{f}(\mathbf{r}_1) \sum_\nu c_{\nu i} \gamma_\nu(\mathbf{r}_1) = \epsilon_i \sum_\nu c_{\nu i} \gamma_\nu(\mathbf{r}_1)$$

$$\sum_\nu c_{\nu i} \int \gamma_\mu^*(\mathbf{r}_1) \hat{f}(\mathbf{r}_1) \gamma_\nu(\mathbf{r}_1) d\mathbf{r}_1 = \epsilon_i \sum_\nu c_{\nu i} \int \gamma_\mu(\mathbf{r}_1) \gamma_\nu(\mathbf{r}_1) d\mathbf{r}_1$$

# Hartree-Fock Theory for $n$ electrons

solving non-linear eigenvalues equations numerically

step 2: expand spatial orbitals in basis functions (basisset)

$$\sum_{\nu} c_{\nu i} \int \gamma_{\mu}^{*}(\mathbf{r}_1) \hat{f}(\mathbf{r}_1) \gamma_{\nu}(\mathbf{r}_1) d\mathbf{r}_1 = \epsilon_i \sum_{\nu} c_{\nu i} \int \gamma_{\mu}(\mathbf{r}_1) \gamma_{\nu}(\mathbf{r}_1) d\mathbf{r}_1$$

express in terms of matrices

$$\sum_{\nu} F_{\mu\nu} c_{\nu i} = \epsilon_i \sum_{\nu} S_{\mu\nu} c_{\nu i}$$

$$\mathbf{F}\mathbf{C} = \mathbf{S}\mathbf{C}\boldsymbol{\epsilon}$$

solution if, and only if

$$|\mathbf{F} - \epsilon_i \mathbf{S}| = 0$$

# Hartree-Fock Theory for $n$ electrons

solving non-linear eigenvalues equations numerically

non-linear:  $\mathbf{F}$  depends on  $\mathbf{C}$

$$F_{\mu\nu} = \int \gamma_\mu^*(\mathbf{r}_1) \hat{h}^0(\mathbf{r}_1) \gamma_\nu(\mathbf{r}_1) d\mathbf{r}_1$$

$$+ 2 \sum_a \int \int \gamma_\mu^*(\mathbf{r}_1) \phi_a^*(\mathbf{r}_2) \frac{1}{r_{12}} \phi_a(\mathbf{r}_2) \gamma_\nu(\mathbf{r}_1) d\mathbf{r}_1 d\mathbf{r}$$

$$- \sum_a \int \int \gamma_\mu^*(\mathbf{r}_1) \phi_a^*(\mathbf{r}_2) \frac{1}{r_{12}} \gamma_\nu(\mathbf{r}_2) \phi_a(\mathbf{r}_1) d\mathbf{r}_1 d\mathbf{r}$$

$$F_{\mu\nu} = h_{\mu\nu}^0$$

$$+ 2 \sum_a \sum_\kappa \sum_\lambda c_{\kappa a}^* c_{\lambda a}^* \int \int \gamma_\mu^*(\mathbf{r}_1) \gamma_\kappa^*(\mathbf{r}_2) \frac{1}{r_{12}} \gamma_\lambda(\mathbf{r}_2) \gamma_\nu(\mathbf{r}_1) d\mathbf{r}_1 d\mathbf{r}$$

$$- \sum_a \sum_\kappa \sum_\lambda c_{\kappa a}^* c_{\lambda a}^* \int \int \gamma_\mu^*(\mathbf{r}_1) \gamma_\kappa^*(\mathbf{r}_2) \frac{1}{r_{12}} \gamma_\nu(\mathbf{r}_2) \gamma_\lambda(\mathbf{r}_1) d\mathbf{r}_1 d\mathbf{r}$$

# Hartree-Fock

## Roothaan-Hall equations

non-linear eigenvalue problem

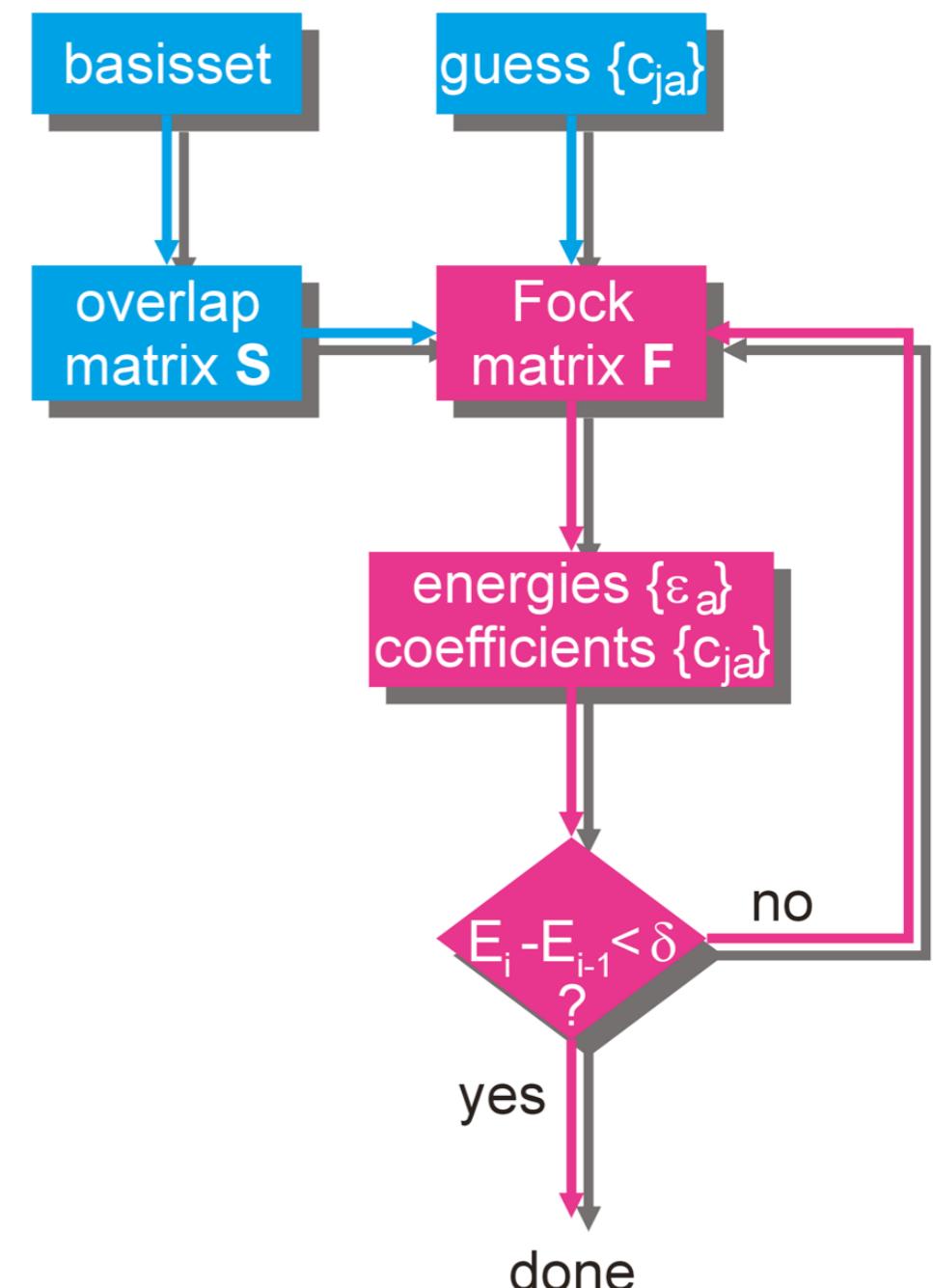
$$Fc = \epsilon_i Sc$$

practical algorithm

iterate until self-consistency

pre-compute integrals of basisset

$$S_{\mu\nu} = \langle \gamma_\mu | \gamma_\nu \rangle$$



# Basissets

minimal basis (1 function per shell)

H-He: 1s (1)

Li-Ne: 1s, 2s, 2p<sub>x</sub>, 2p<sub>y</sub>, 2p<sub>z</sub> (5)

Na-Ar: 1s, 2s, 2p<sub>x</sub>, 2p<sub>y</sub>, 2p<sub>z</sub>, 3s, 3p<sub>x</sub>, 3p<sub>y</sub>, 3p<sub>z</sub> (9)

Slater-type orbitals

computationally demanding

$$f_{1s}(\zeta, \mathbf{r}) = \exp[-\zeta r]$$

Gaussian-type orbitals

computationally convenient

$$g_{1s}(\alpha, \mathbf{r}) = (8\alpha^3/\pi^3)^{1/4} \exp[-\alpha r^2]$$

$$g_{2p_x}(\alpha, \mathbf{r}) = (128\alpha^5/\pi^3)^{1/4} x \exp[-\alpha r^2]$$

$$g_{3d_{xy}}(\alpha, \mathbf{r}) = (2048\alpha^7/\pi^3)^{1/4} xy \exp[-\alpha r^2]$$

# Basissets

## Gaussian-type orbitals

computationally convenient, but not as accurate as Slater-type orbitals

linear combination (contraction) of several gaussians (primitives)

### STO-3G

$$\phi_{1s}^{\text{CGF}} = \sum_i^3 d_{i,1s} g_{1s}(\alpha_{i,1s})$$

$$\phi_{2s}^{\text{CGF}} = \sum_i^3 d_{i,2s} g_{1s}(\alpha_{i,2sp})$$

$$\phi_{2p}^{\text{CGF}} = \sum_i^3 d_{i,2p} g_{2p}(\alpha_{i,2sp})$$

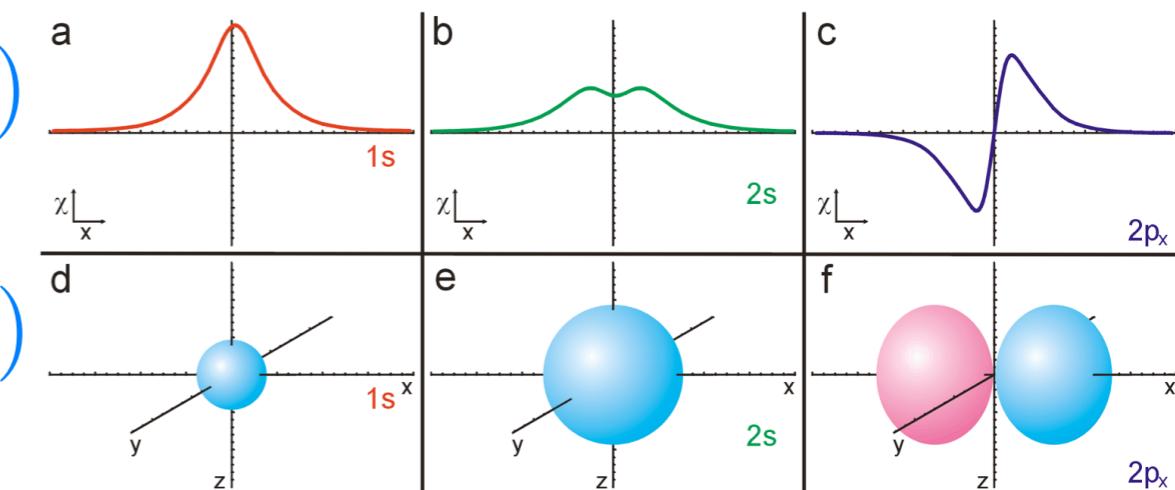
least-square fit to Slater orbitals

$$\min \left[ \int (\phi_{1s}^{\text{SF}}(\mathbf{r}) - \phi_{1s}^{\text{CGF}}(\mathbf{r})) d\mathbf{r} \right]^2$$

$$\min \left[ \int (\phi_{2s}^{\text{SF}}(\mathbf{r}) - \phi_{2s}^{\text{CGF}}(\mathbf{r})) d\mathbf{r} - \int (\phi_{2p}^{\text{SF}}(\mathbf{r}) - \phi_{2p}^{\text{CGF}}(\mathbf{r})) d\mathbf{r} \right]^2$$

notation (for 1st row elements)

$$(6s3p/3s)[2s1p/1s]$$



# Basissets

## Double-Zeta basis

two basisfunctions (contractions) per valence orbital

3-21G

$$(6s3p/3s)[3s2p/2s]$$

H-He:

$$\phi'_{1s} = \sum_i^2 d_{i,1s} g_{1s}(\alpha'_{i,1s})$$

$$\phi''_{1s} = g_{1s}(\alpha''_{i,1s})$$

Li-Ne:

$$\phi_{1s} = \sum_i^3 d_{i,1s} g_{1s}(\alpha_{i,1s})$$

$$\phi'_{2s} = \sum_i^2 d'_{i,2s} g_{1s}(\alpha'_{i,2sp})$$

$$\phi''_{2s} = g_{1s}(\alpha''_{i,2sp})$$

$$\phi'_{2p} = \sum_i^2 d'_{i,2p} g_{2p}(\alpha'_{i,2sp})$$

$$\phi''_{2p} = g_{2p}(\alpha''_{i,2sp})$$

# Basissets

## Double-Zeta basis

two basisfunctions (contractions) per valence orbital

4-31G

$$(8s4p/4s)[3s2p/2s]$$

H-He:

$$\phi'_{1s} = \sum_i^3 d_{i,1s} g_{1s}(\alpha'_{i,1s})$$

$$\phi''_{1s} = g_{1s}(\alpha''_{i,1s})$$

Li-Ne:

$$\phi_{1s} = \sum_i^4 d_{i,1s} g_{1s}(\alpha_{i,1s})$$

$$\phi'_{2s} = \sum_i^3 d'_{i,2s} g_{1s}(\alpha'_{i,2sp})$$

$$\phi''_{2s} = g_{1s}(\alpha''_{i,2sp})$$

$$\phi'_{2p} = \sum_i^3 d'_{i,2p} g_{2p}(\alpha'_{i,2sp})$$

$$\phi''_{2p} = g_{2p}(\alpha''_{i,2sp})$$

# Basissets

## Double-Zeta basis

two basisfunctions (contractions) per valence orbital

6-31G

$$(10s4p/4s)[3s2p/2s]$$

H-He:

$$\phi'_{1s} = \sum_i^3 d_{i,1s} g_{1s}(\alpha'_{i,1s})$$

$$\phi''_{1s} = g_{1s}(\alpha''_{i,1s})$$

Li-Ne:

$$\phi_{1s} = \sum_i^6 d_{i,1s} g_{1s}(\alpha_{i,1s})$$

$$\phi'_{2s} = \sum_i^3 d'_{i,2s} g_{1s}(\alpha'_{i,2sp})$$

$$\phi''_{2s} = g_{1s}(\alpha''_{i,2sp})$$

$$\phi'_{2p} = \sum_i^3 d'_{i,2p} g_{2p}(\alpha'_{i,2sp})$$

$$\phi''_{2p} = g_{2p}(\alpha''_{i,2sp})$$

# Basissets

## Double-Zeta basis with polarization functions

two basisfunctions (contractions) per valence orbital

Li-Ne: 3d functions (\*)

H-He: 2p functions (\*\*)

3-21G\*, 4-31G\*, 6-31G\*, 6-31G\*\*

note 6 d-functions (i.e. 5 3d functions + 3s)

$3d_{xx}, 3d_{yy}, 3d_{zz}, 3d_{xy}, 3d_{yz}, 3d_{zx}$

linear combinations of 5 pure 3d functions + 3s:

$3d_{xy}, 3d_{x^2-y^2}, 3d_{yz}, 3d_{zx}, 3d_{z^2}, 3s_{x^2+y^2+z^2}$

contractions

6-31G\*:  $(11s4p1d/4s)[4s2p1d/2s]$

6-31G\*\*:  $(11s4p1d/4s1p)[4s2p1d/2s1p]$

# Molecular Quantum Mechanics

## molecular orbitals

linear combination of atomic orbitals (hydrogen)

$$\phi_i(r) = \sum_j c_{ji} \chi_j(r)$$

example: hydrogen molecule, minimal basis

$$\phi_1(r) = \chi_1(r) + \chi_2(r)$$

$$\phi_2(r) = \chi_1(r) - \chi_2(r)$$



# Molecular Quantum Mechanics

## molecular orbitals

linear combination of atomic orbitals (hydrogen)

$$\phi_i(r) = \sum_j c_{ji} \chi_j(r)$$

example: hydrogen molecule, minimal basis

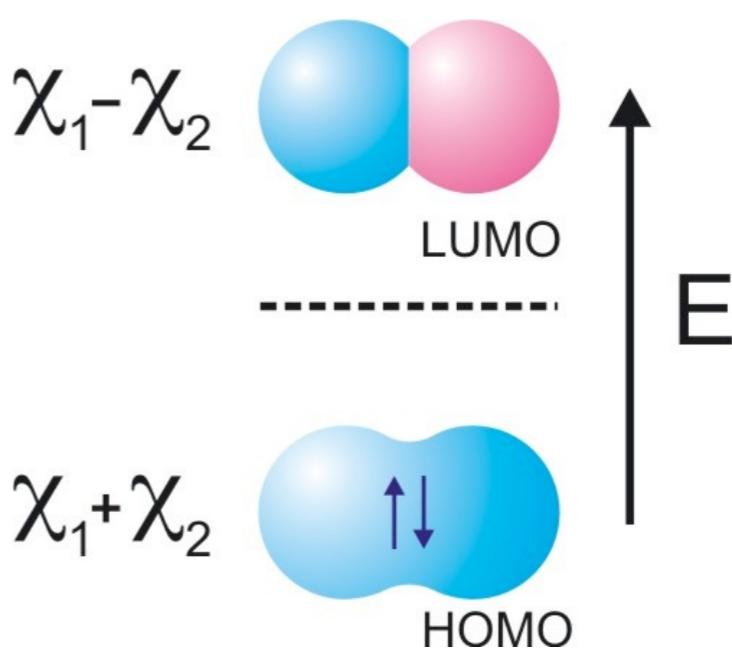
$$\phi_1(r) = \chi_1(r) + \chi_2(r)$$

$$\Psi(r_1, r_2) = |\phi_1(r_1)\overline{\phi_1}(r_2)|$$

$$\phi_2(r) = \chi_1(r) - \chi_2(r)$$



dissociation!!!



# Molecular Quantum Mechanics

## limitations of Hartree-Fock theory

- ▶ mean-field: no electron correlation

## beyond Hartree-Fock

- ▶ configuration interaction (ci)

$$\begin{aligned}\Psi = & C_0 |\phi_a(r_1)\phi_b(r_2) \dots \phi_k(r_{n-1})\phi_l(r_{n-1})| + \\ & C_1 |\phi_k(r_1)\phi_b(r_2) \dots \phi_k(r_{n-1})\phi_l(r_{n-1})| + \\ & C_2 |\phi_k(r_1)\phi_l(r_2) \dots \phi_k(r_{n-1})\phi_l(r_{n-1})| + \dots\end{aligned}$$

- ▶ multiconfiguration self-consistent field (mcscf)

# Electron correlation

## definition

$$E^{\text{corr.}} = E^{\text{exact}} - E^{\text{HF}}$$

## Configuration Interaction

sum over all possible Slater determinants/spin adapted-configuration state functions

general expression for correlation energy (not so useful...)

truncated CI

MCSCF/CASSCF

size consistency

## Perturbation theory

Slater determinants, single, double, triple, ... excitations

## Combinations of CI and PT

MCSCF with  $x$ th order perturbation ( $x=2$ )

Concludes the ab initio, next is density functional theory